

Cooling and Seebeck effects in Double-Barrier Semiconductor Heterostructures

M. Bescond^{(1,2,3)*}

¹ IM2NP/CNRS, UMR 7334, Aix-Marseille University, France

² Institute of Industrial Science and INQIE, The University of Tokyo, Japan

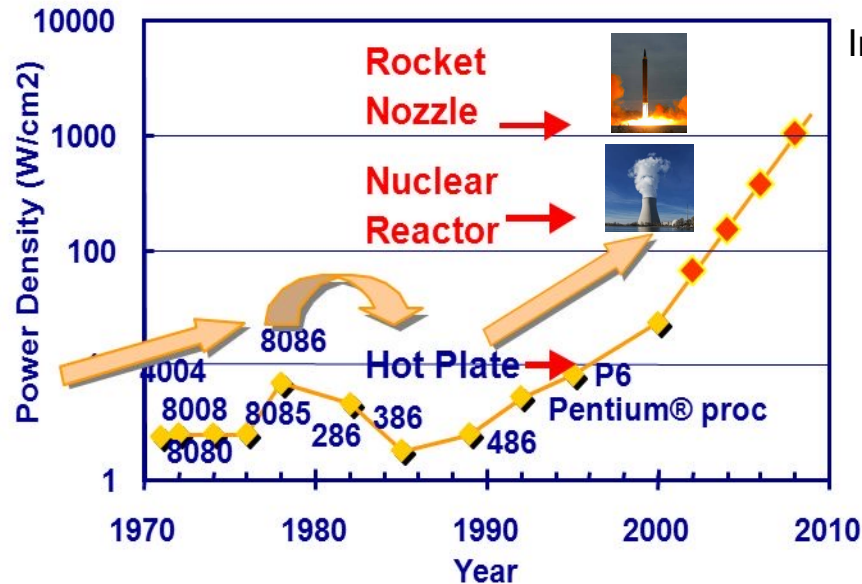
³ Before at LIMMS/CNRS-IIS, IRL 2820, Tokyo, Japan

Thematic day « Beyond Fourier »

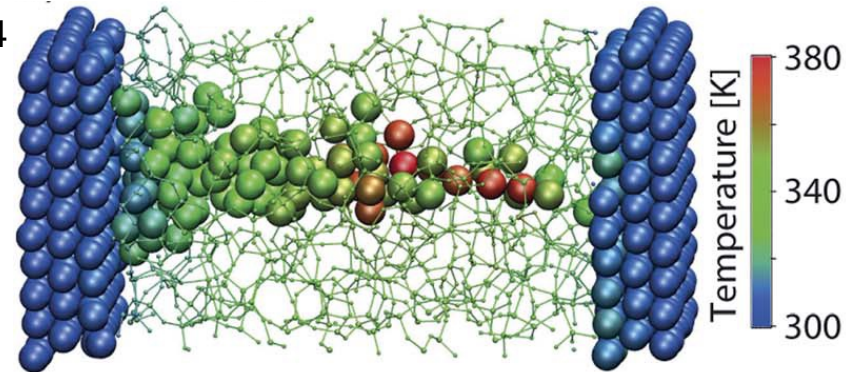
The 9th of September 2022

Cooling at the nanoscale

➤ Self-heating: scientific and industrial issues



Intel, 2004



CBRAM

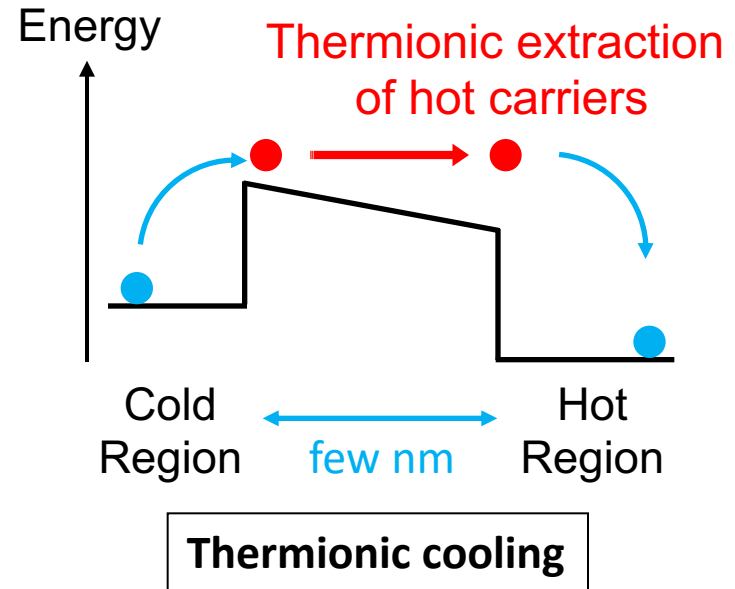
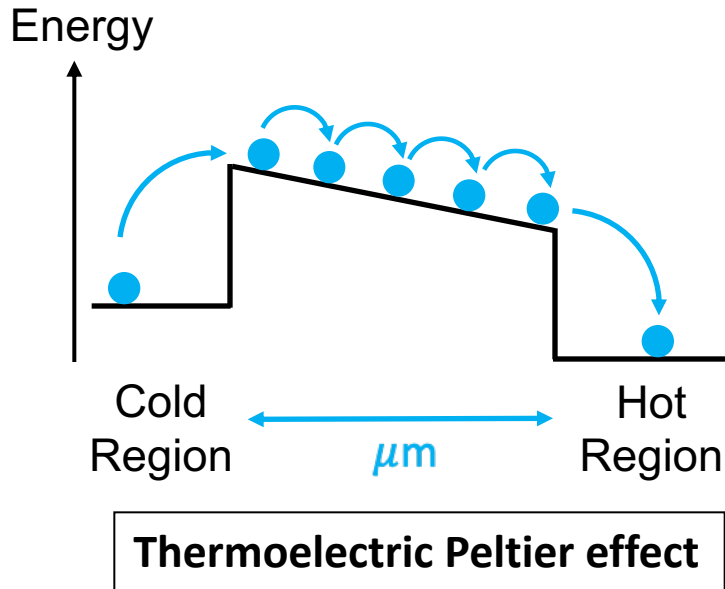
Nanoscale Adv., 2, 2648 (2020)

M. Luisier, ETH Zurich

- Significant reduction of lifetimes and performances.
- “Bulk” refrigeration is extremely power consuming.

Urgent need of local source of cooling

Thermionic cooling*



- Devices working in non-equilibrium regime.

 **Non-equilibrium \rightarrow Highest cooling power!**

- Exploratory field: **strong theoretical support.**

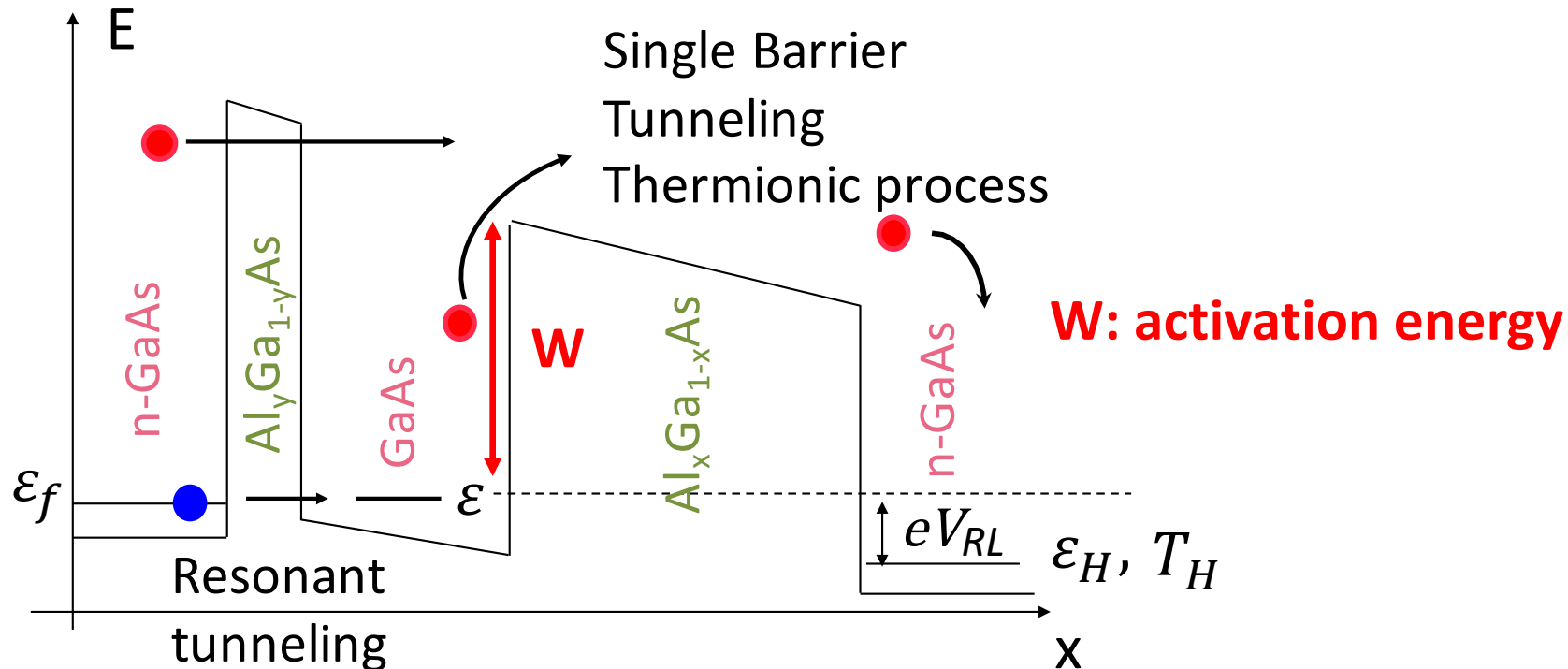
Goal: use nano-structures to improve cooling efficiency.

*G. D. Mahan, *J. Appl. Phys.*, **76**, 4362 (1994).

*G. D. Mahan and L. M. Woods, *Phys. Rev. Lett.*, **80**, 4016 (1998).

Original thermionic cooling devices

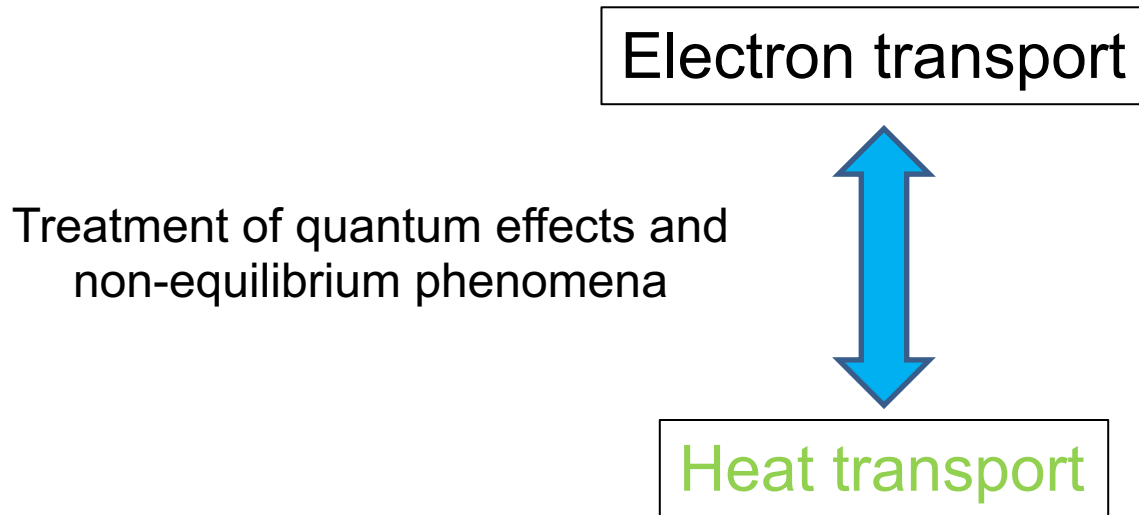
- Coupling localized state and tunneling barrier*:



- Main idea: injecting cold electrons and extracting hot electrons.
- Investigate the concept of local temperature at the nanoscale.

Simulations Requirements

Coupling of electron and heat transport



- Electron → Non-Equilibrium Green's Function (NEGF) 🤔
 - Effective mass approach, k.p, tight-binding, ab initio...
- Heat → Non-Equilibrium Green's Function (NEGF) 😊

Much less documented topic:

Harmonic (ballistic): N. Mingo and L. Yang, *Phys. Rev. B* **68**, 245406 (2003).

NEGF for phonon

- NEGF for electrons: Schrödinger equation

$$(E\mathbf{I} - \mathbf{H})\bar{\psi} = \bar{0} \quad \longrightarrow \quad [G^R(E)]$$

- NEGF for phonons: dynamical equation

$$\omega^2 \mathbf{M} + \Phi_{nm}^{ij} \bar{\mathbf{R}} = \bar{0} \quad \text{with} \quad \Phi_{nm}^{ij} = \frac{\partial^2 V^{harm}}{\partial R_n^i \partial R_m^j}$$

Vibration frequency

Mass of the atoms (matrix)

2nd order FC

Displacement of the atoms

Retarded Green's function

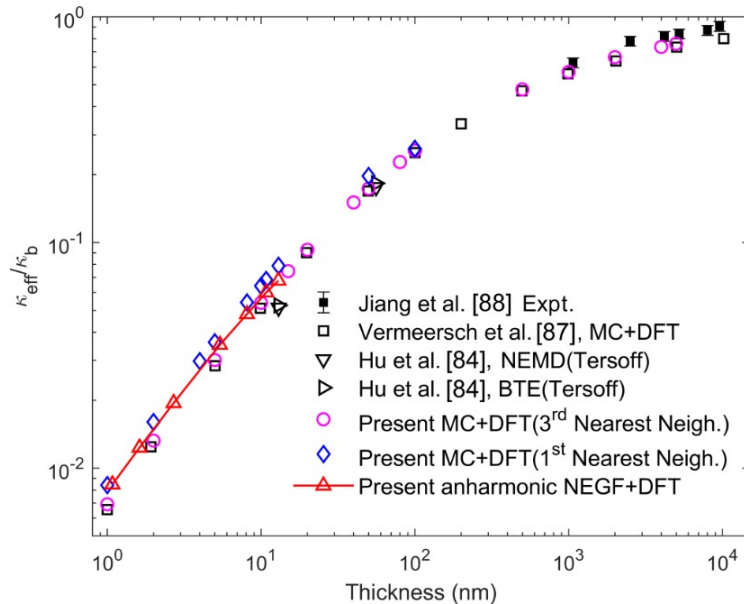
$$\mathbf{G}^R(\omega) = [\omega^2 \mathbf{I} - \Phi - \Sigma^R(\omega)]^{-1}$$

- Anharmonic: N. Mingo, *Phys. Rev. B* **74**, 125402 (2006). ([junctions](#))
- J. S. Wang, N. Zeng, J. Wang, and C. K. Gan, *Phys. Rev. E* **75**, 061128 (2007).
- M. Luisier, *Phys. Rev. B* **86**, 245407 (2012). ([nanowires](#))
- K. Miao *et al.*, *Appl. Phys. Lett.* **108**, 113107 (2016). ([Büttiker Probes](#))
- J. H. Dai and Z. T. Tian, *Phys. Rev. B* **101**, 041301 (2020). ([interfaces](#))
- R. Rhyner and M. Luisier, *Phys. Rev. B* **89**, 235311 (2014). ([NEGF coupling e-/ph!!](#))

NEGF for phonon

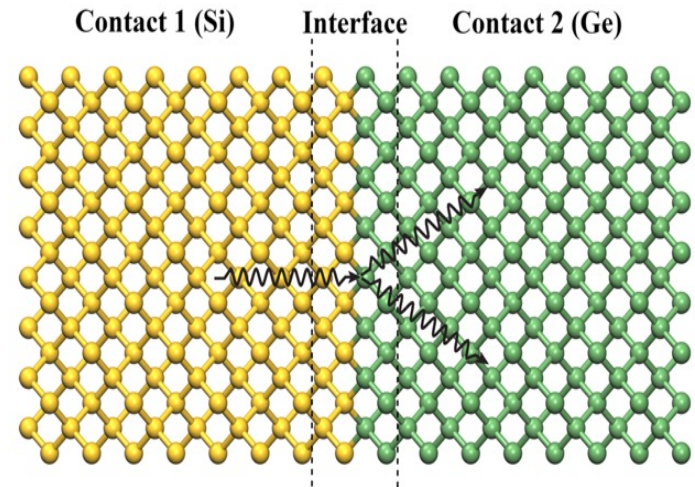
➤ Anharmonic phonon NEGF development: additional issues

Thermal conductivity Si film



Y. Guo, et al., *Phys. Rev. B*, **102**, 195412 (2020).

Si/Ge interface



Y. Guo, et al., *Phys. Rev. B* **103** (17), 174306 (2021).

➤ Anharmonicity:

- 1) Non-local treatment of the phonon-phonon self-energies!
- 2) 4th order force constant might be needed...

→ NEGF for phonons can be only applied to rather small (tens of nanometers) systems.

Alternative approach:

NEGF for electrons

+

Heat equation



M. Bescond *et al.* *J. Phys.: Condens. Matter* **30**, 064005 (2018).

NEGF + Heat equation

➤ Non-equilibrium Green's function coupled to heat equation*

NEGF equations for electrons

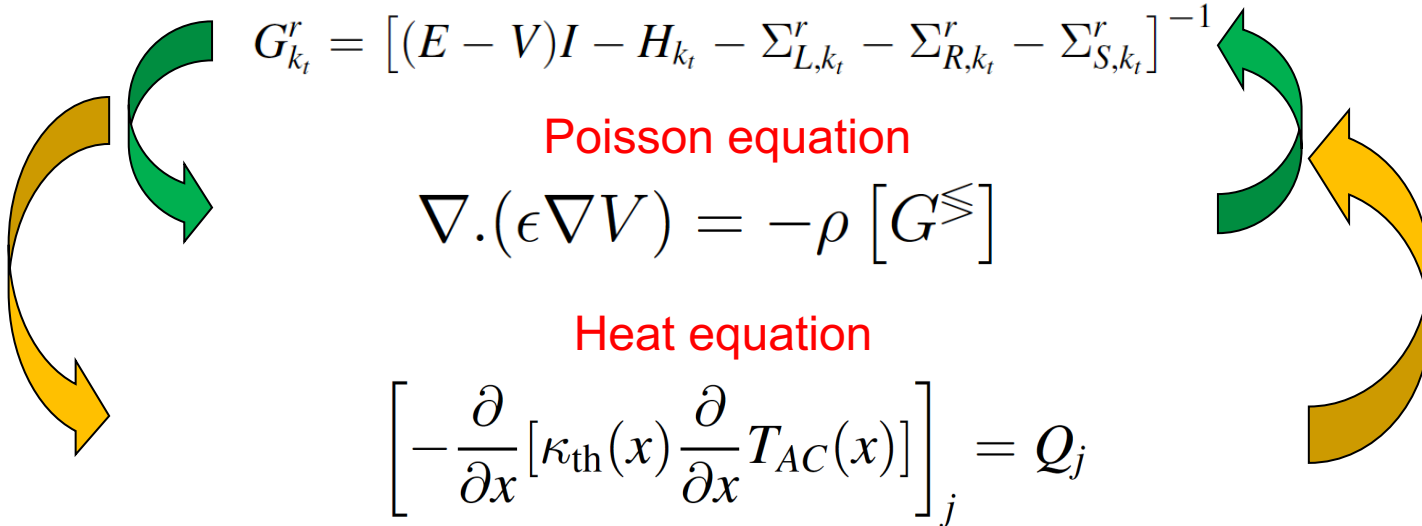
$$G_{k_t}^r = [(E - V)I - H_{k_t} - \Sigma_{L,k_t}^r - \Sigma_{R,k_t}^r - \Sigma_{S,k_t}^r]^{-1}$$

Poisson equation

$$\nabla \cdot (\epsilon \nabla V) = -\rho [G^{\lessgtr}]$$

Heat equation

$$\left[-\frac{\partial}{\partial x} \left[\kappa_{\text{th}}(x) \frac{\partial}{\partial x} T_{AC}(x) \right] \right]_j = Q_j$$

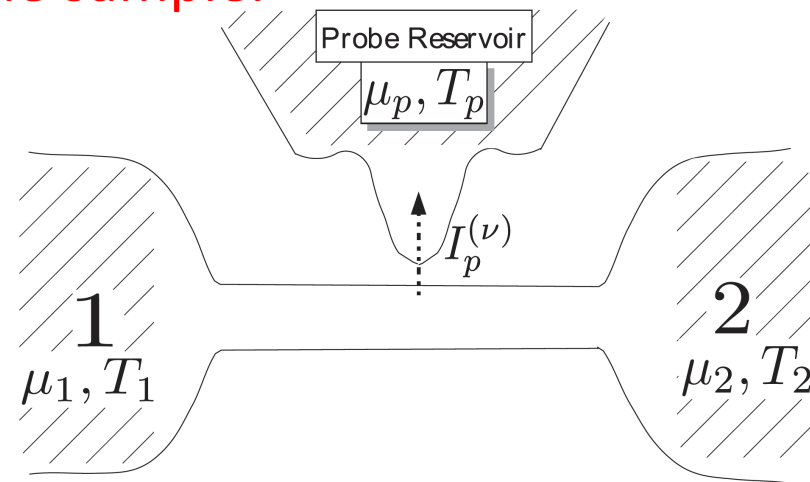


- Most of physical properties: current, electron density, LDOS, local phonon temperatures, cooling power, efficiency...
- But... We are in a strong non-equilibrium regime... $T_{AC} \neq T_{POP} \neq T_e$
- Temperature of electrons in the active region???

Electronic temperature: virtual probe technique

- System out of equilibrium:
Electronic and lattice temperatures usually not coincide.
- Accurate electronic temperature measurement (i.e. that follows the thermodynamic laws) requires simultaneously local voltage measurement.^{1,2}
- Technique: vanish net charge current ($I_p^{(0)}$) **and** net heat current ($I_p^{(1)}$) into the probe.
--> probe in local equilibrium with the sample.

$$I_p^{(\nu)} = 0, \quad \nu \in \{0, 1\}$$

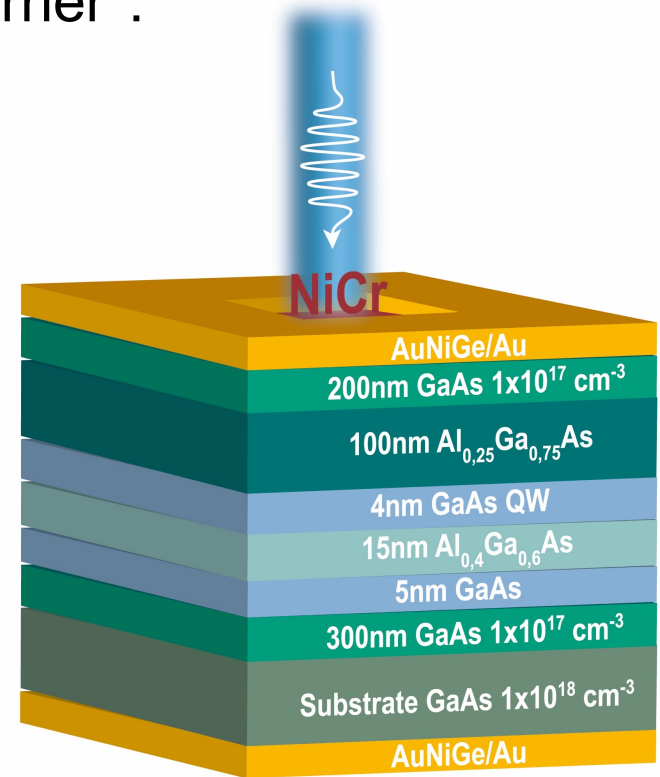
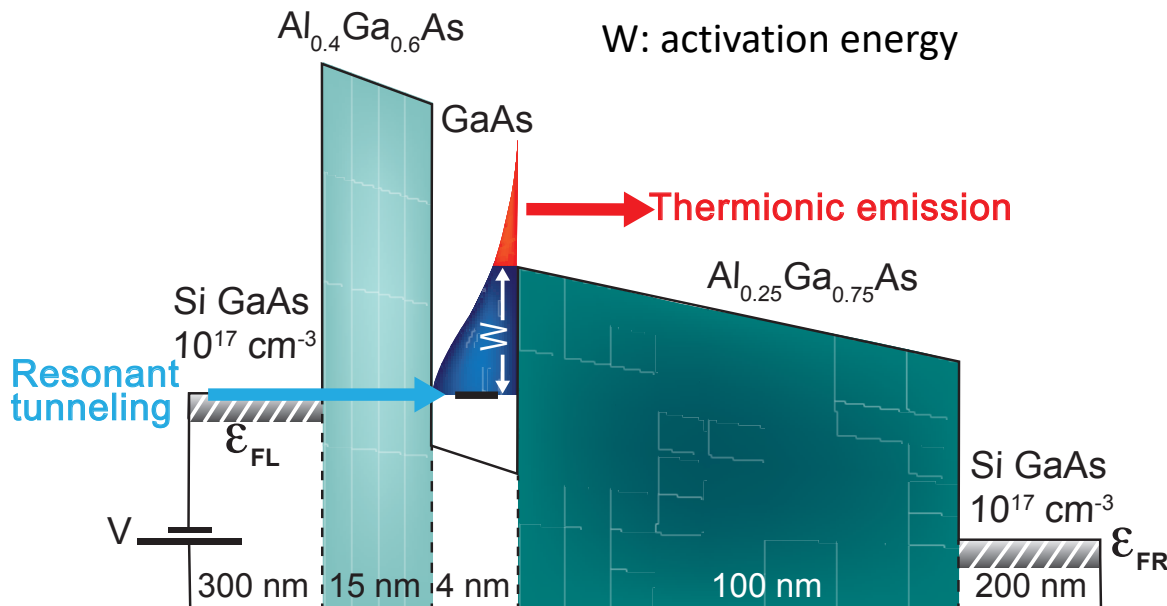


¹C. A. Stafford, *Phys. Rev. B* **93**, 245403 (2016).

²A. Shastry and C. A. Stafford, *Phys. Rev. B* **94**, 155433 (2016).

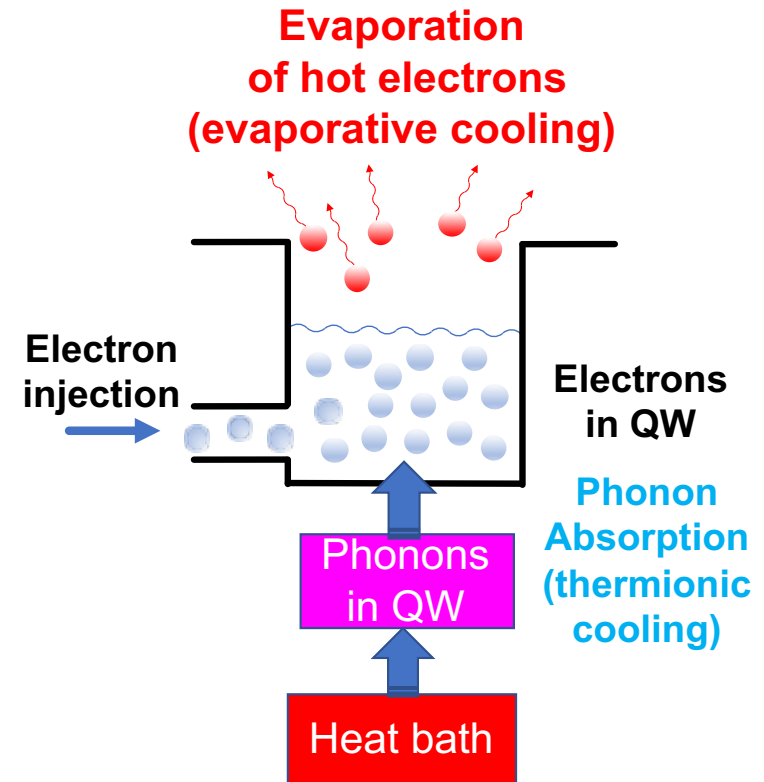
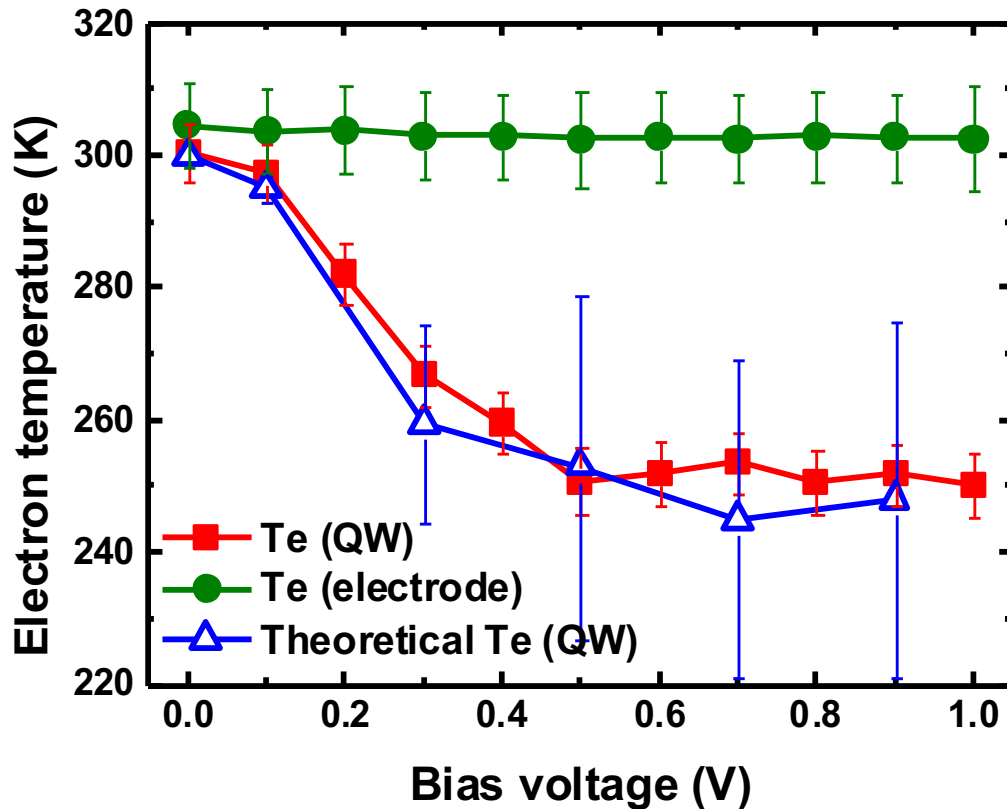
Experimental implementation

- Coupling localized state and tunneling barrier*:



- Sample fabrication: Molecular beam epitaxy (MBE).
- Temperature of electron T_e in the quantum well (QW).

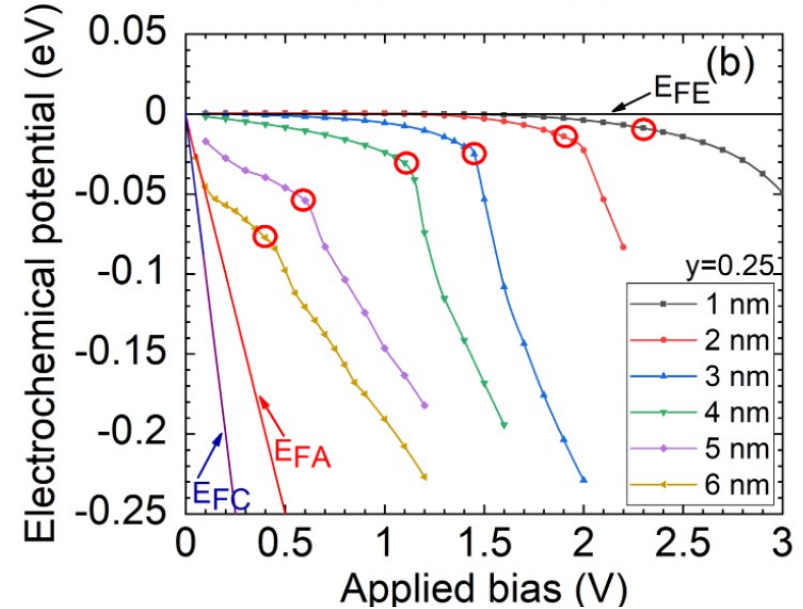
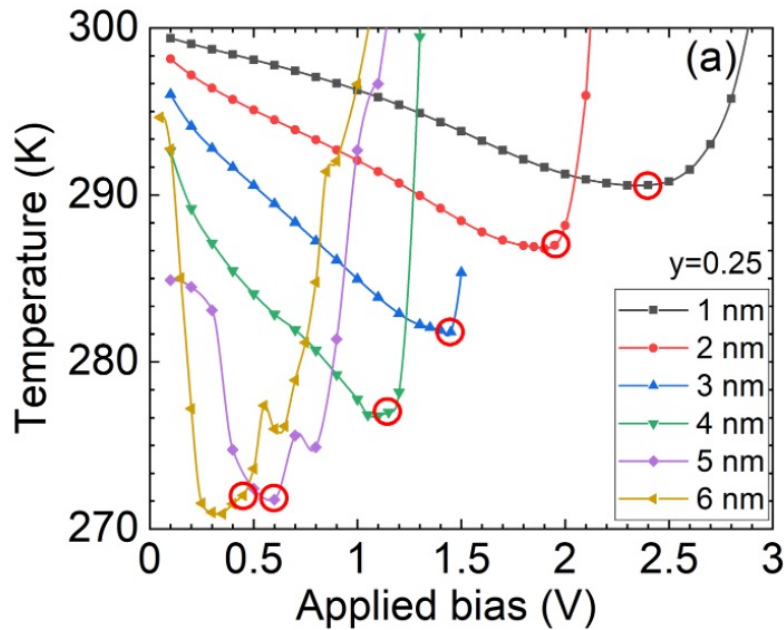
Electron Temperature(s)*



- T_e in electrodes constant.
- T_e in the QW decreases by 50K due to evaporative cooling.

*A. Yangui, M. Bescond, T. Yan, N. Nagai, and K. Hirakawa, *Nature Commun.* **10**, 4504 (2019).

Maximum electron cooling*



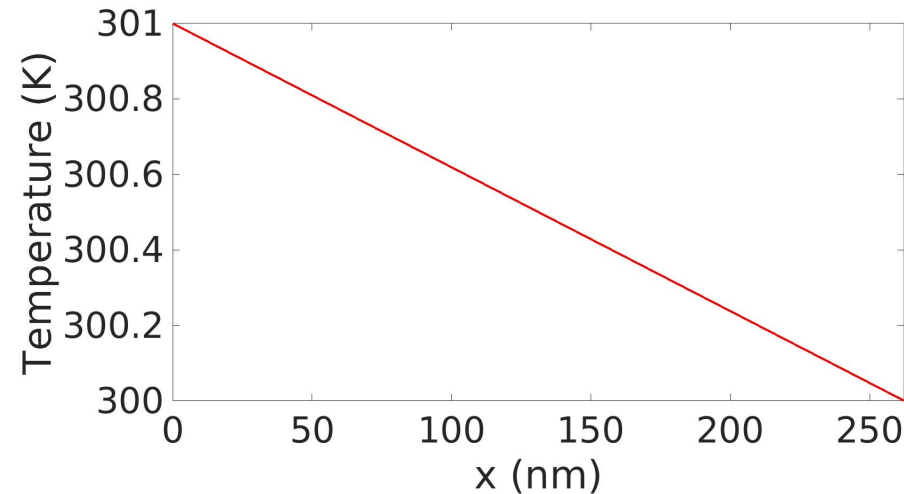
- Temperature minimum coincides very well with the resonance condition (highest current \rightarrow best energy filtering).
- Temperature reduction increases with L_{Emit} .
- Electrochemical potential: best cooling when $R_{\text{Emit}} \approx R_{\text{Coll}}$.

*M. Bescond, et al., *Phys. Rev. Appl.* **17**, 014001 (2022).

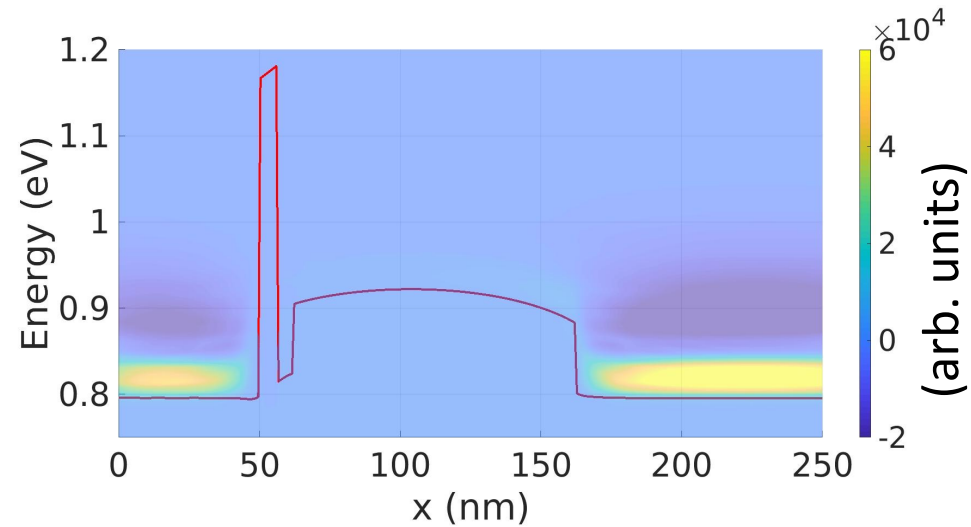
Thermoelectric properties

Lattice temperature gradient

$$\Delta T = 1 \text{ K}$$



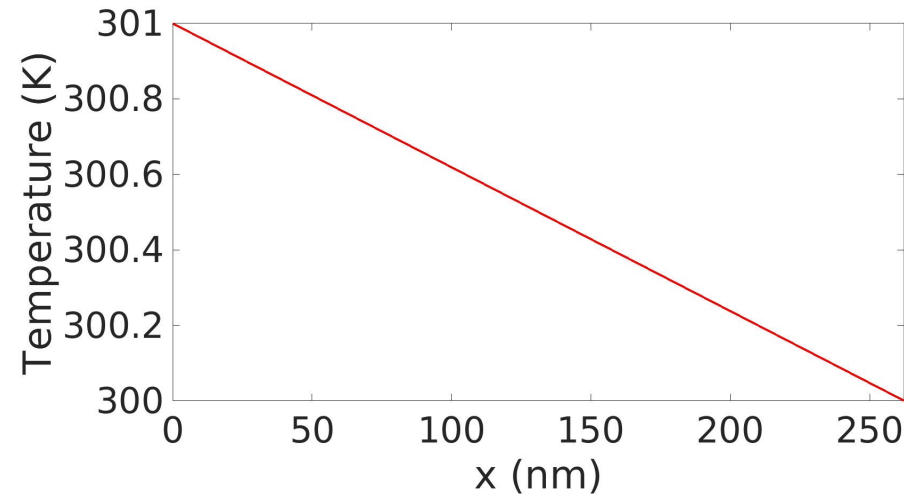
Current spectrum



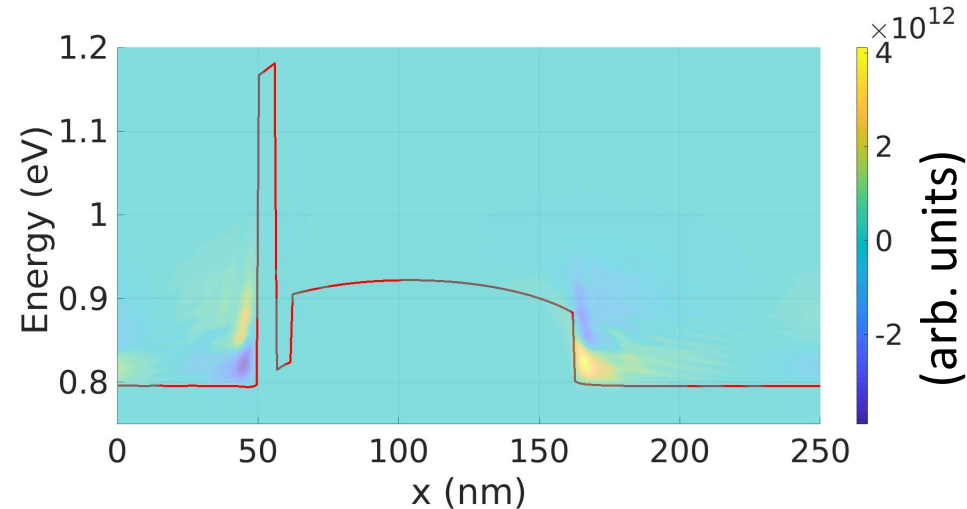
- Two current components in the access regions.
- At the barrier: phonon emission/absorption and back flow towards the contacts

Lattice temperature gradient

$$\Delta T = 1 \text{ K}$$



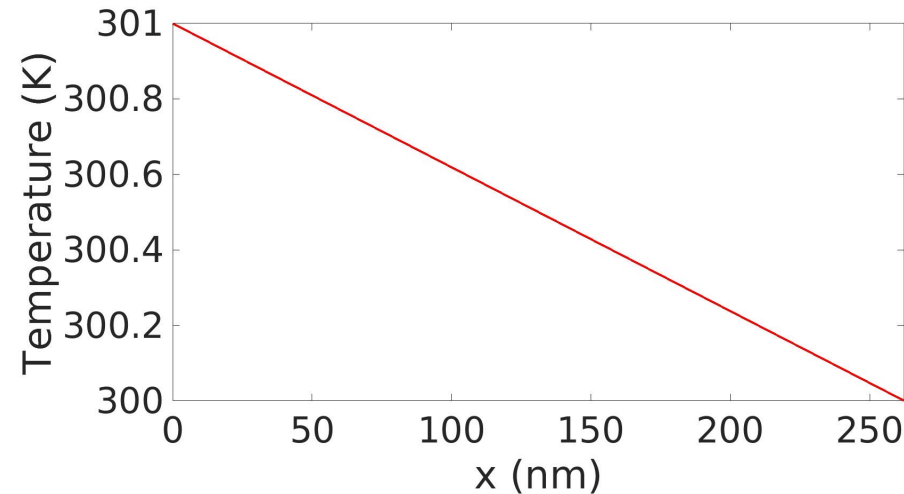
Derivative of energy current



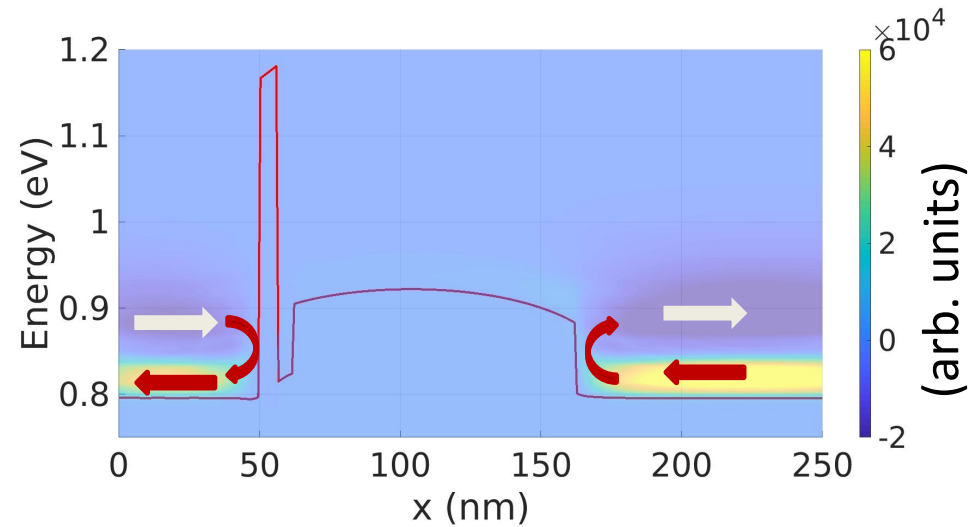
- Two current components in the access regions.
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Lattice temperature gradient

$$\Delta T = 1 \text{ K}$$



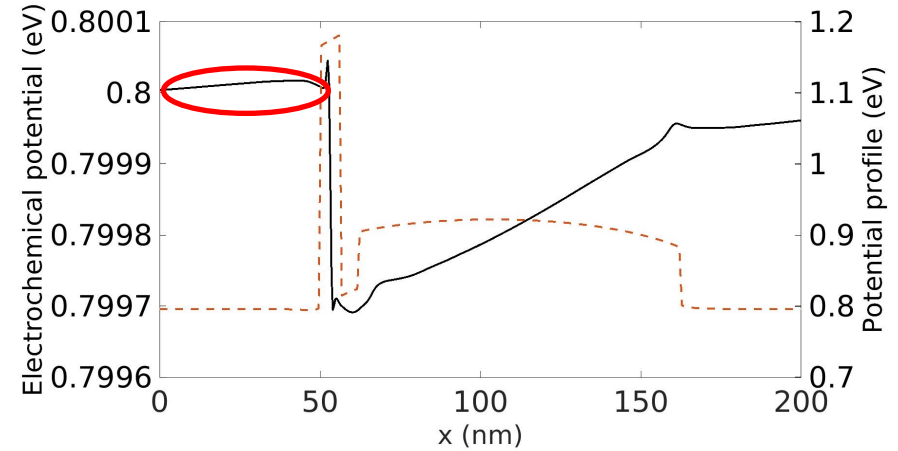
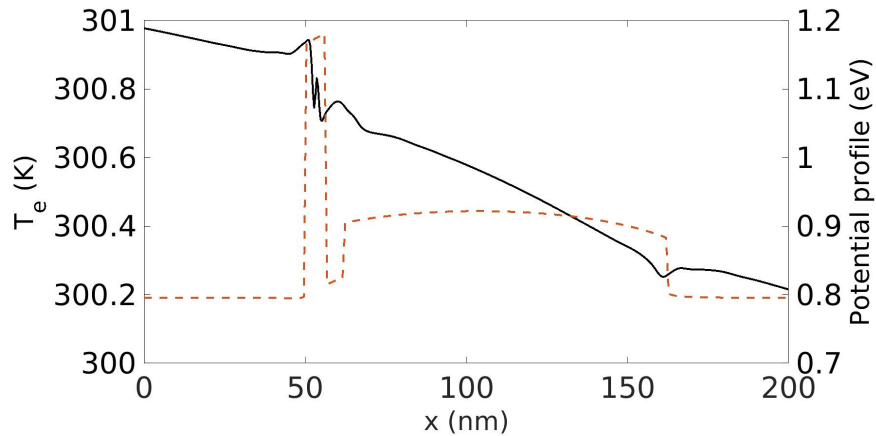
Current spectrum



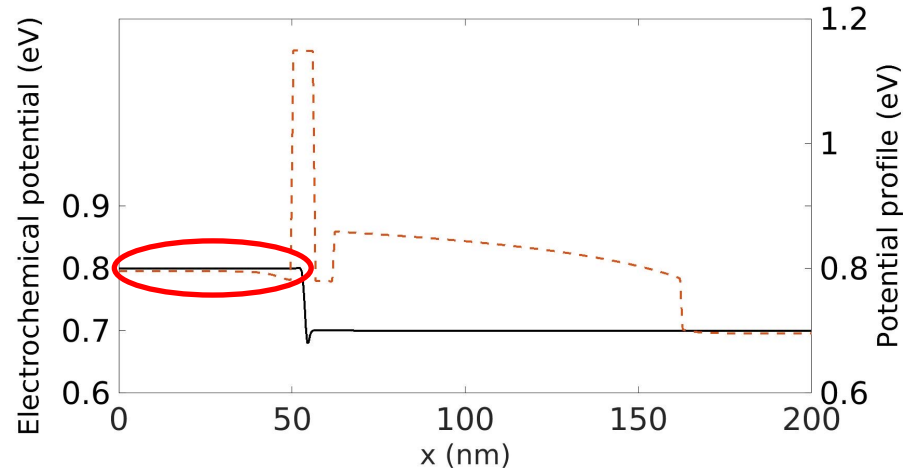
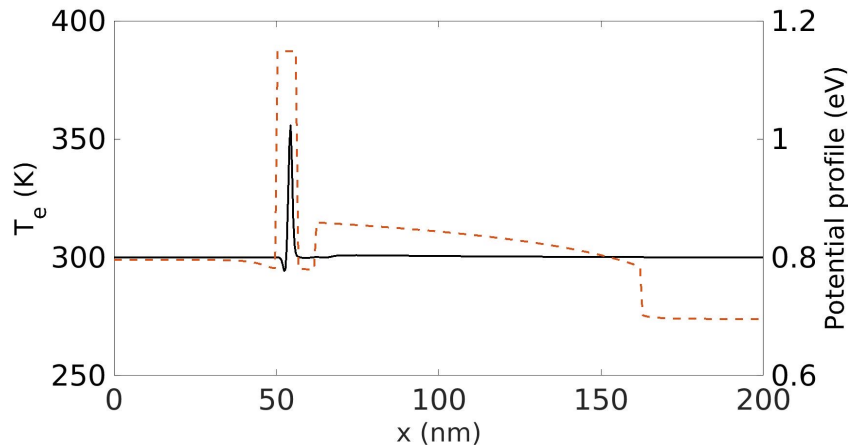
- Two current components in the access regions.
- At the barrier: phonon emission/absorption and back flow towards the contacts

Electron temperature and Chemical potential

$$\Delta T = 1 \text{ K}, \Delta V = 0 \text{ V}$$



$$\Delta T = 0 \text{ K}, \Delta V = 0.1 \text{ V}$$

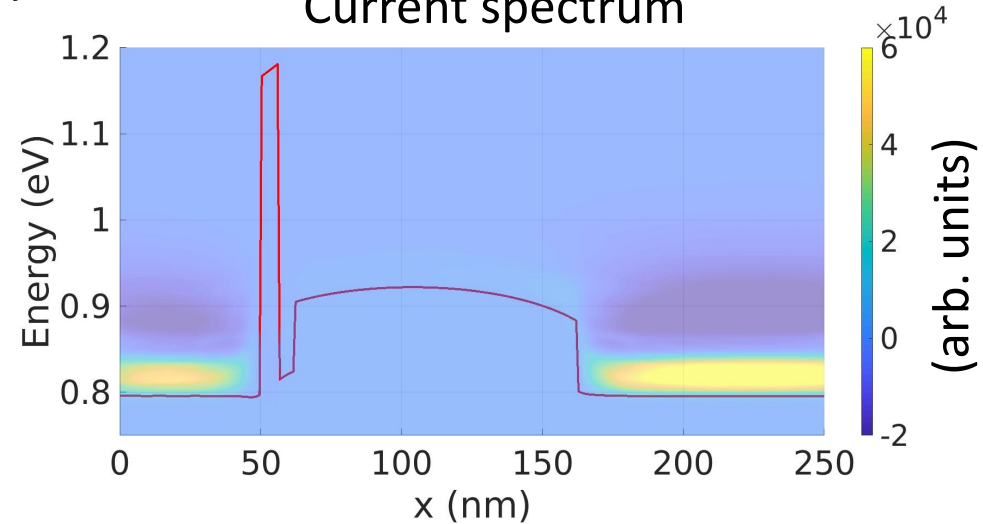
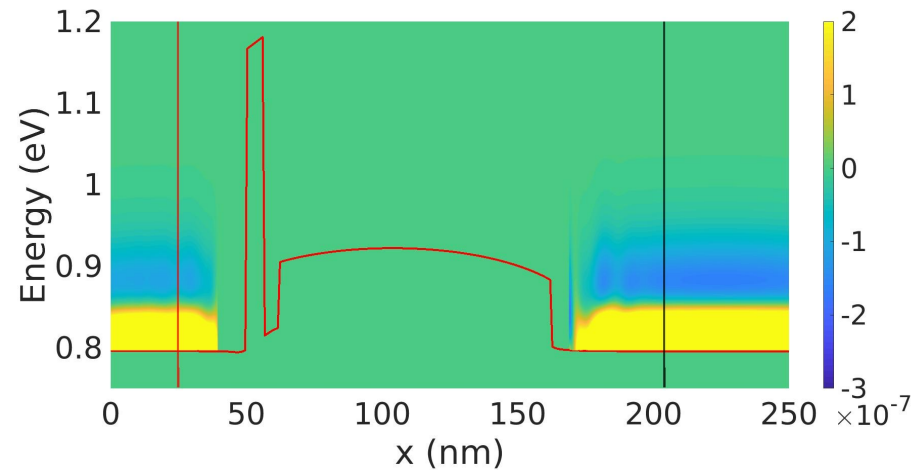


Difference of the Fermi-Dirac distributions

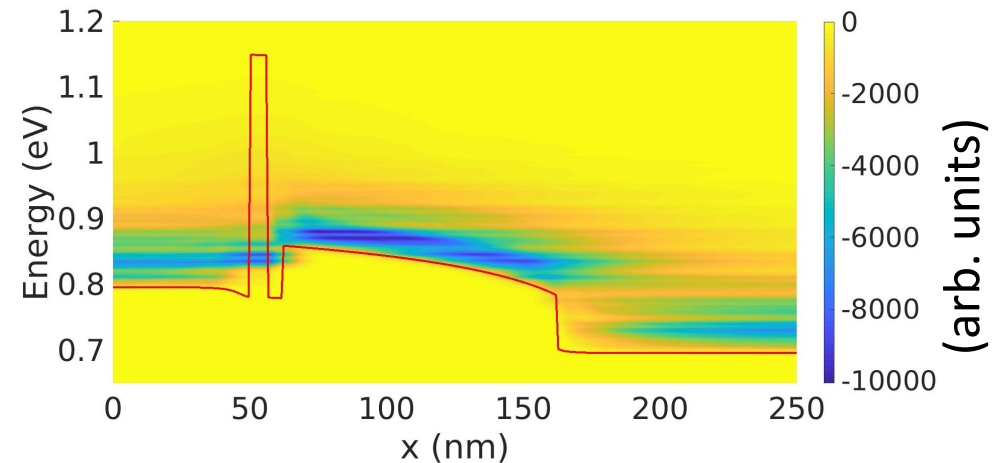
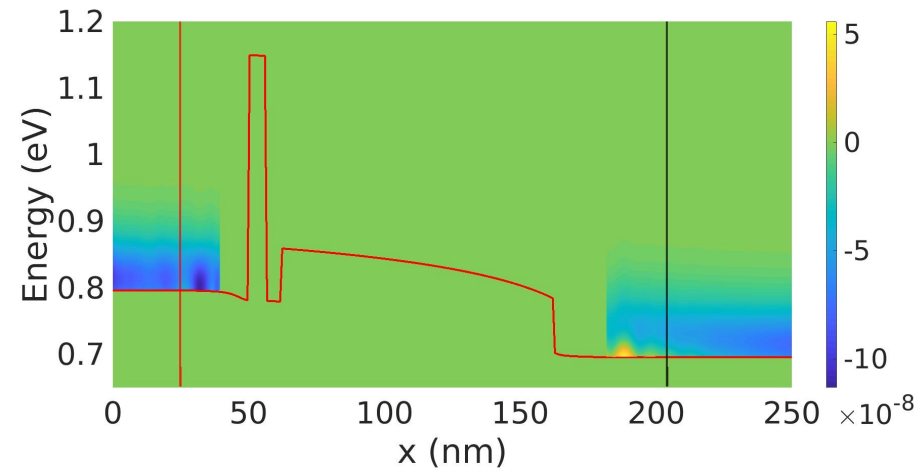
$\Delta T = 1$ K, $\Delta V = 0$ V

Δf_{FD}

Current spectrum



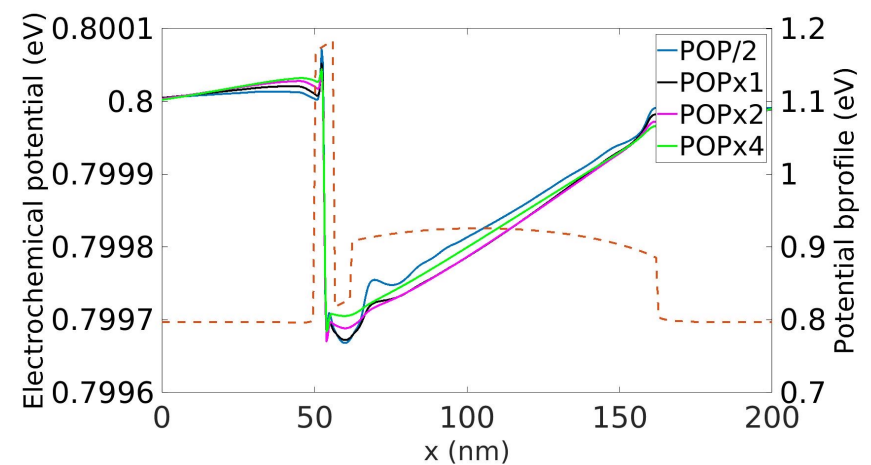
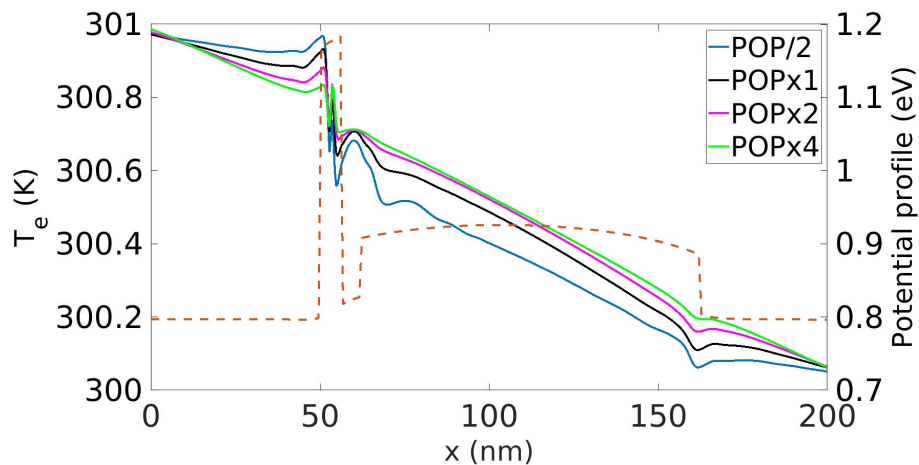
$\Delta T = 0$ K, $\Delta V = 0.1$ V



Origin of the reverse current component

- Reverse current: due to the variation of μ in the device.
- Increase of μ required to maintain the electron density with a ΔT_e .

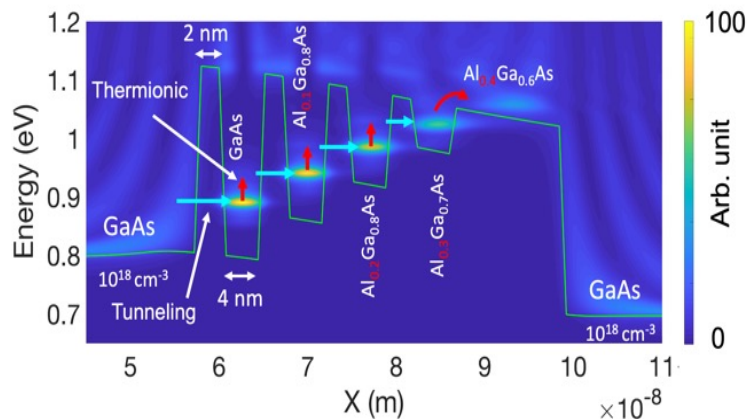
Impact of the electron-phonon coupling



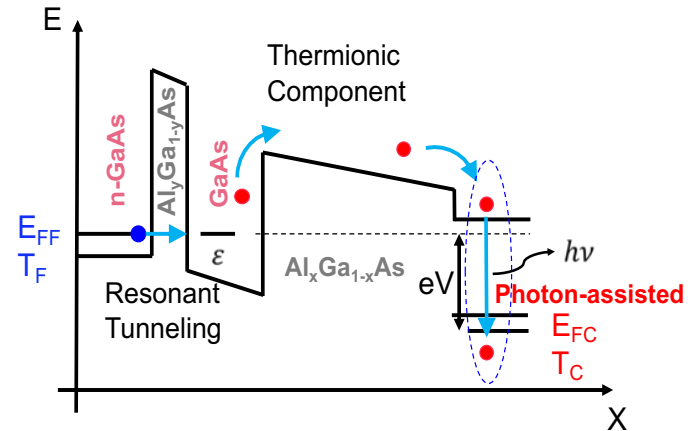
- Increasing coupling: T_e follow T_{Lattice}
 - ➔ Larger T_e decrease
 - ➔ Larger μ increase

Conclusion

- **Coupling NEGF and heat equation:** good insight of the physical properties of the thermoelectric systems (good agreement with experiment).
- **Conception and optimization of devices**



Quantum cascade cooler



Opto-thermionic pumping



Novel generation of cooling devices

ANR GELATO (Oct. 2021- Apr.2025)

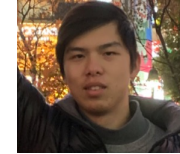
LIMMS, CNRS, IIS-University of Tokyo:



➤ Pr. Kazuhiko Hirakawa



Dr. Chloé Salhani
Post-doc



Mr. Xiangyu Zhu
PhD student

NQS group, IM2NP, CNRS, Aix-Marseille University:



Pr. N. Cavassilas



Dr. F. Micheline
(Assoc. Pr.)



Dr. A.-M. Daré
(Assoc. Pr.)



Prof. M. Lannoo
(Emeritus)



Celine Belabbas (Intern)



Dr. A. Crépieux
(Assoc. Pr. at CPT)



Mr. G. Etesse
(PhD student)

Thank you!