

Thermal transfer at the Atomic scale in Amorphous materials and Nanostructured amorphous samples

Anne Tanguy

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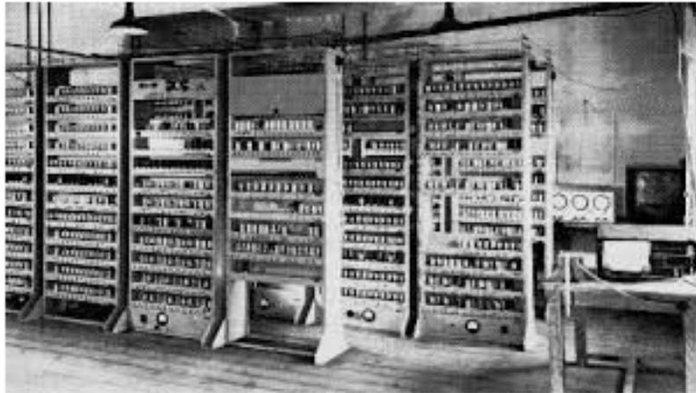
&

LaMCoS, UMR CNRS5259 / INSA-LYON

20 Avenue Albert Einstein
69 621 Villeurbanne, France

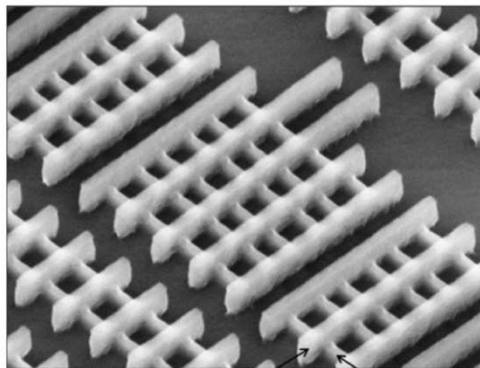


It is important to convert or to evacuate (to manage) the heat at small scales



1950's

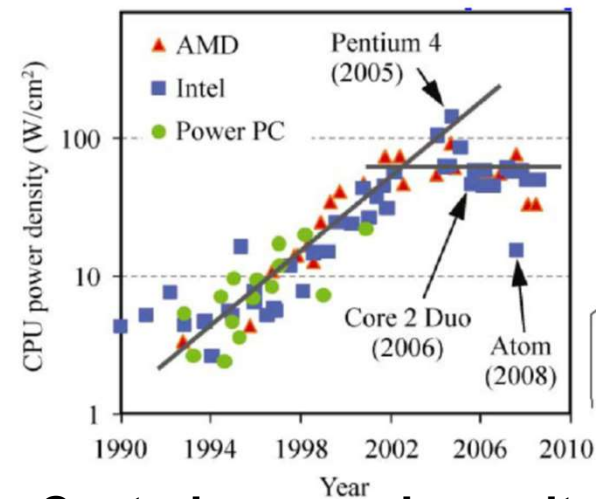
22 nm Tri-Gate Transistor (Intel, 2011)



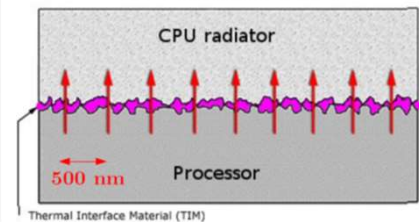
Gates Fins



2000's

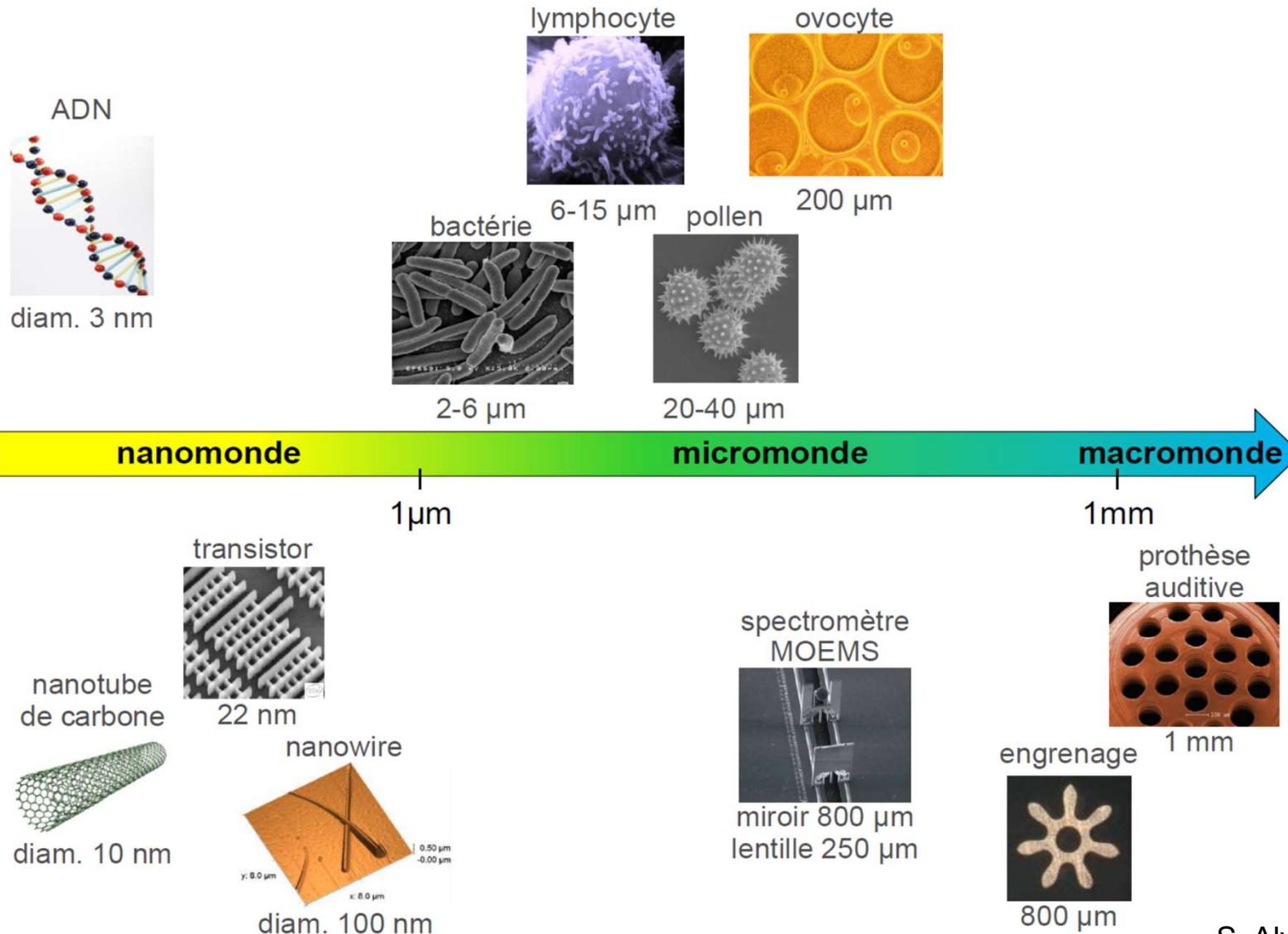


Contact using a TIM



Central processing unit power dissipation

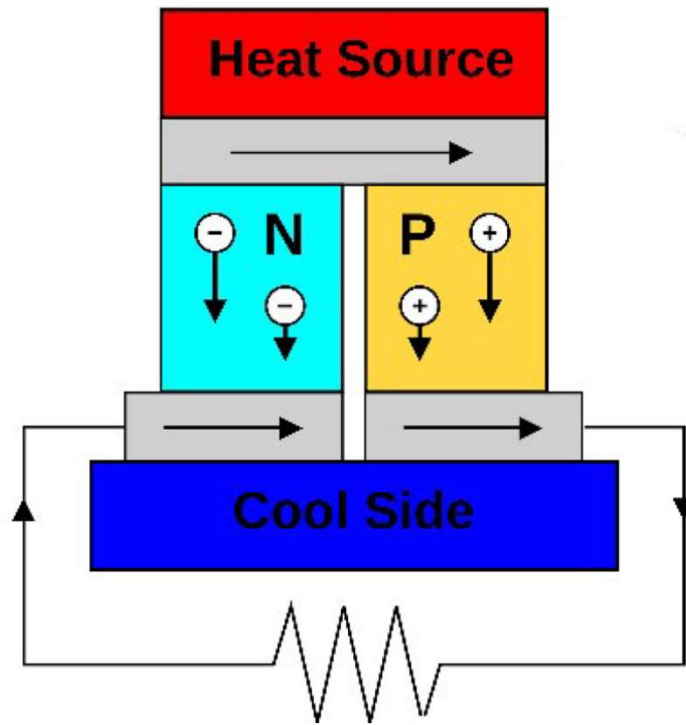
Laboratoire de Mécanique des Contacts et des Structures



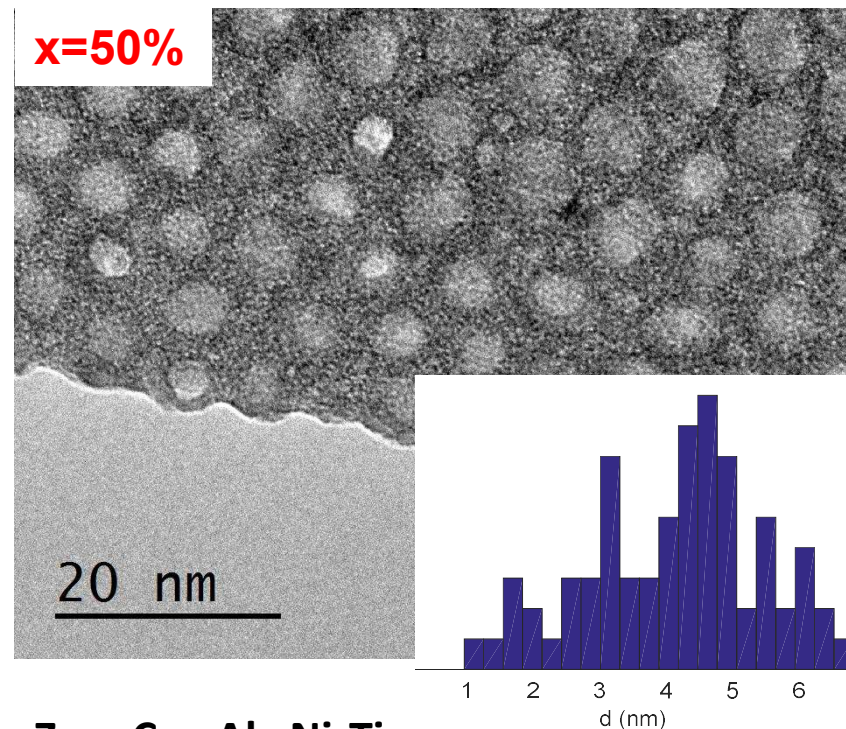
Other example: Energy harvesting, Thermoelectric Efficiency: Merit Factor

$$Z = \frac{S^2 \sigma}{\lambda_{ph} + \lambda_e}$$

σ — electric conductivity
 $\lambda_{ph} + \lambda_e$ — thermal conductivity



Property	Metals	Semiconductors	Insulators
S (μVK^{-1})	~ 5	~ 200	~ 1000
σ ($\Omega^{-1}\text{cm}^{-1}$)	$\sim 10^6$	$\sim 10^3$	$\sim 10^{-12}$
Z (K^{-1})	$\sim 3 \times 10^{-6}$	$\sim 2 \times 10^{-3}$	$\sim 5 \times 10^{-17}$



$\text{Zr}_{52.5}\text{Cu}_{27}\text{Al}_{10}\text{Ni}_8\text{Ti}_{2.5}$

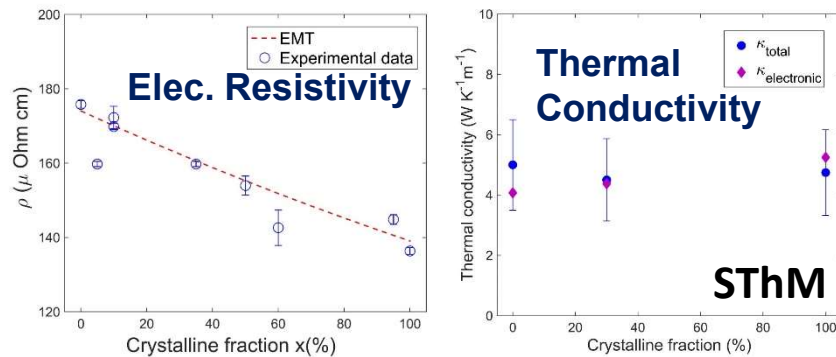
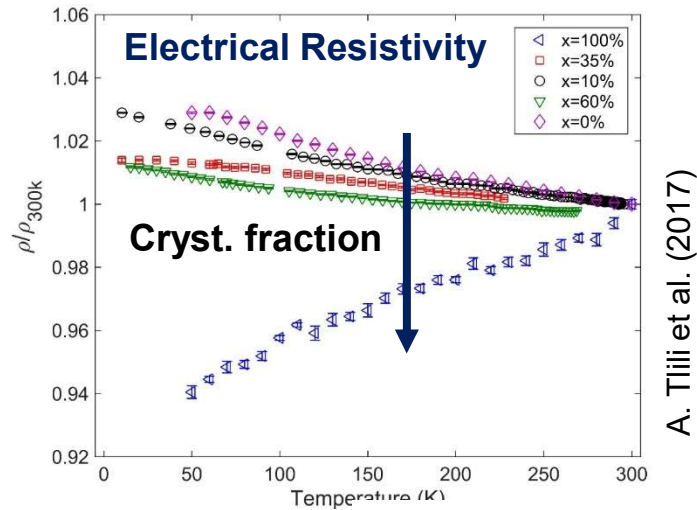
Crystallites mean diameter $\sim 4 \pm 2$ nm

V. Giordano et al. (2017)

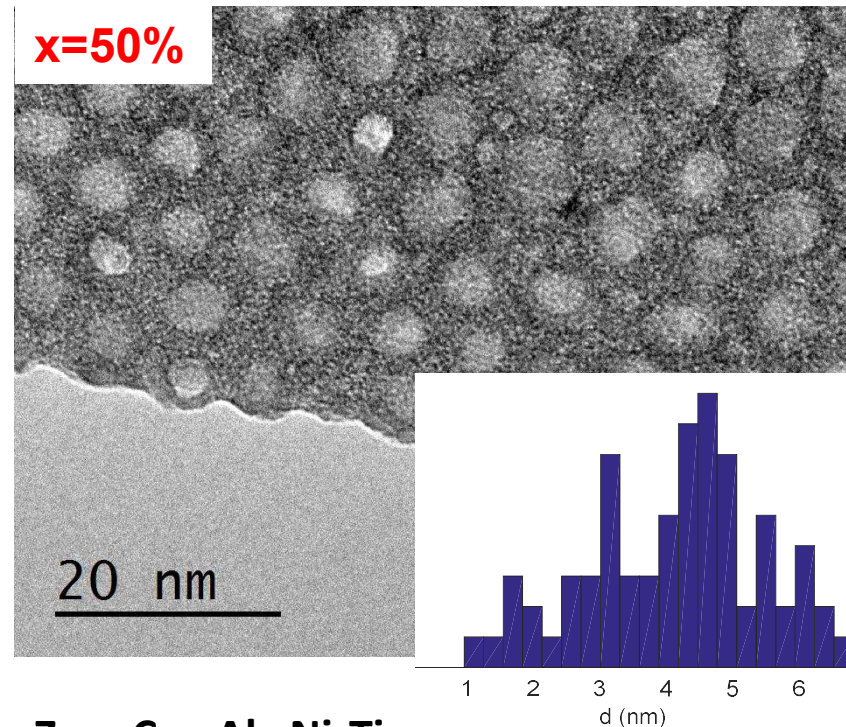
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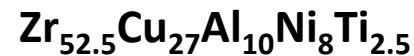
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V. Giordano et al. (2017)



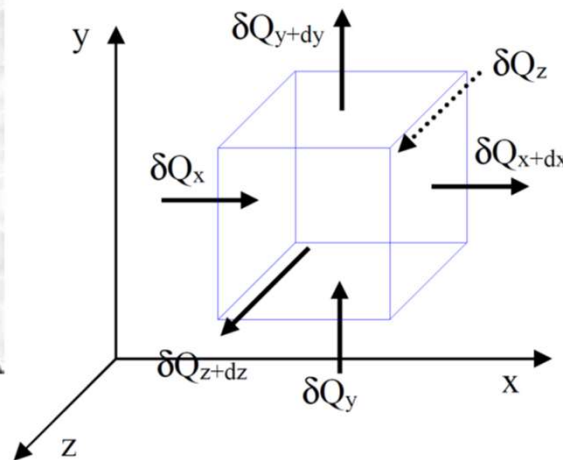
Crystallites mean diameter $\sim 4 \pm 2$ nm

Understanding the Thermal Conductivity at small scales

Out-of-Equilibrium Energy Fluxes:



J.-B. J. Fourier
(1768-1830)



Flux of heat

Fourier's Law

$$\delta Q_x = -k_x \frac{\partial T}{\partial x} dy dz dt$$

Thermal Conductivity

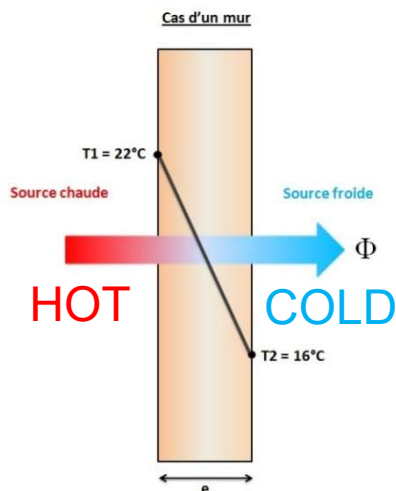
Energy Conservation:

$$\delta Q_x + \delta Q_y + \delta Q_z + q dV dt = \delta Q_{x+dx} + \delta Q_{y+dy} + \delta Q_{z+dz} + \rho C \cdot dV \cdot \frac{\partial T}{\partial t} dt$$

$$\frac{\partial}{\partial x} \left(k_x \frac{\partial T}{\partial x} \right) + \frac{\partial}{\partial y} \left(k_y \frac{\partial T}{\partial y} \right) + \frac{\partial}{\partial z} \left(k_z \frac{\partial T}{\partial z} \right) + q = \rho C \frac{\partial T}{\partial t}$$

$$\Delta T + \frac{q}{k} = \frac{1}{D} \frac{\partial T}{\partial t}$$

$$D = \frac{k}{\rho C} \text{ Diffusivity}$$



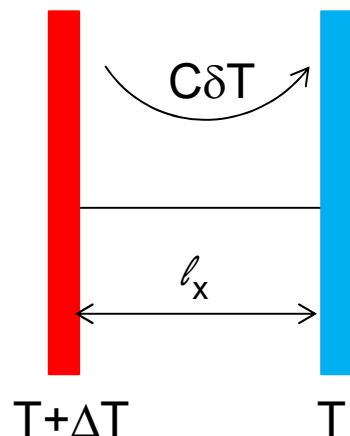
Thermal Conductivity in crystals (Kittel):

Heat carriers = electrons, phonons

In crystals, phonons = quantum of vibrational energy $\left(n + \frac{1}{2}\right)\hbar\omega$

Thermal energy transfer = random diffusive process, « collisions » of heat carriers

Kinetic Theory of Thermal Conductivity:



$$\delta T = \frac{dT}{dx} \ell_x = \frac{dT}{dx} v_x \tau \quad \ell_x, \text{ phonons mean free path}$$

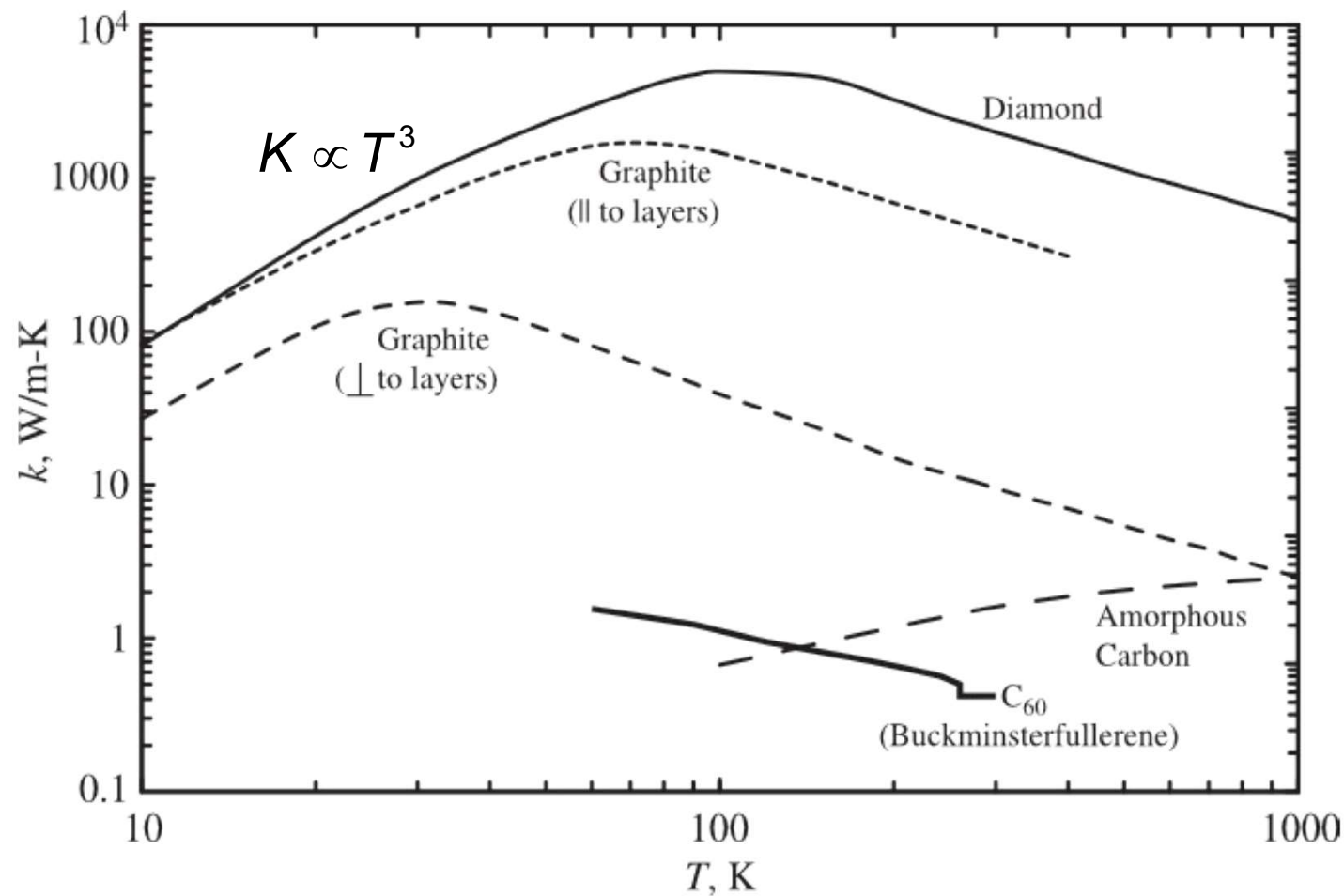
$$\text{Net flux of Energy: } j_q = -\frac{N}{V} \langle v_x C \delta T \rangle = -\frac{N}{V} \left\langle v_x C \frac{dT}{dx} \ell_x \right\rangle \approx -\frac{1}{3} \frac{N}{V} C \|v\| \ell_x \frac{dT}{dx}$$

$$\Rightarrow \kappa \approx \frac{1}{3} C \|v\| \ell_x \quad \text{Heat Conductivity}$$

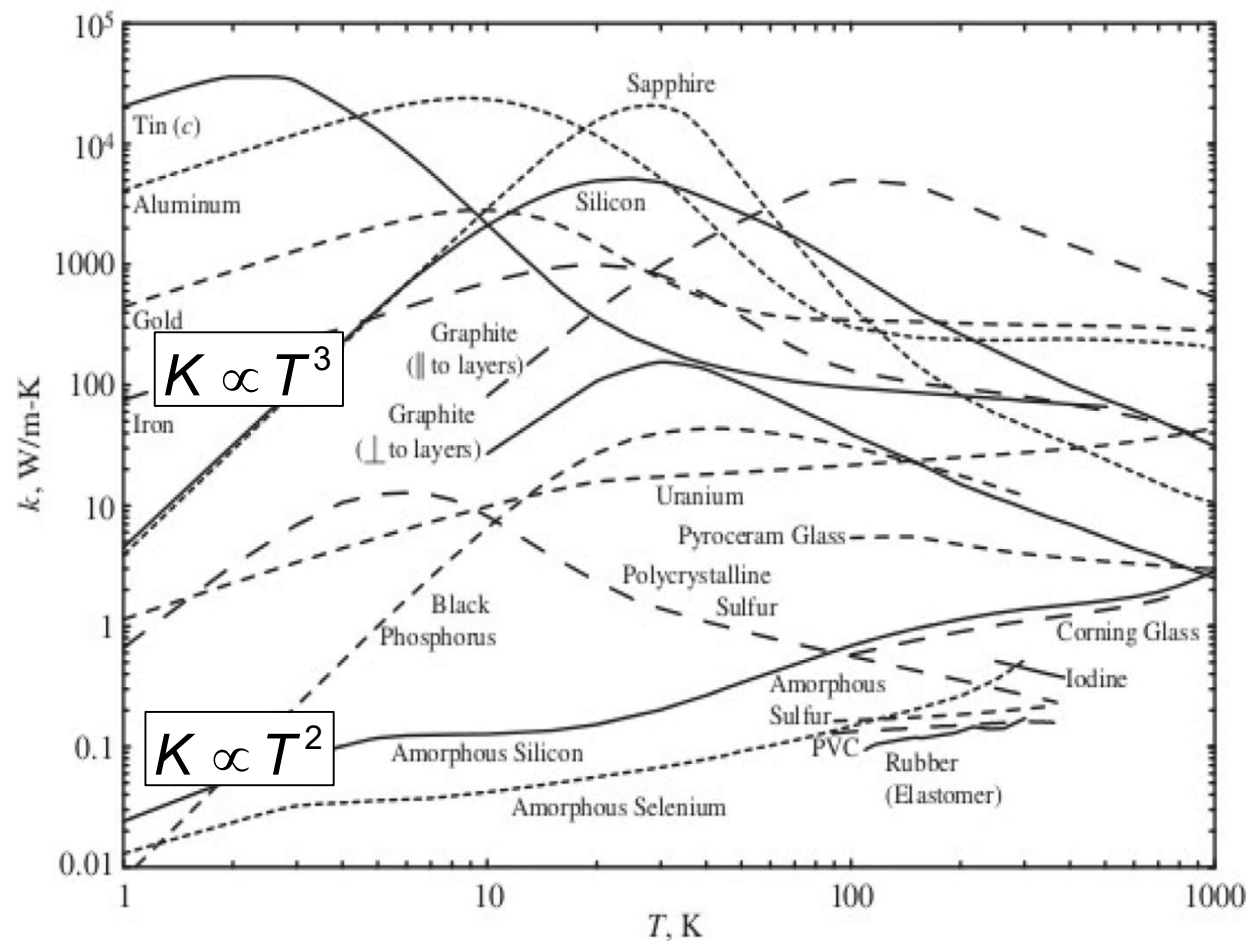
Geometrical scattering: $\ell_x \approx L$ $\kappa \propto C \propto T^3$
crystal boundary, lattice imperfections...

Scattering by other phonons: $\ell_x \propto \frac{1}{n_{\text{phonons}}} \propto \frac{1}{T}$
anharmonicity, umklapp processes...

Some Measurements of Thermal Conductivity:



Some Measurements of Thermal Conductivity:



Thermal Conductivity

plateau:

Resonant scattering
of phonons by
« quasi-local vibrations »?

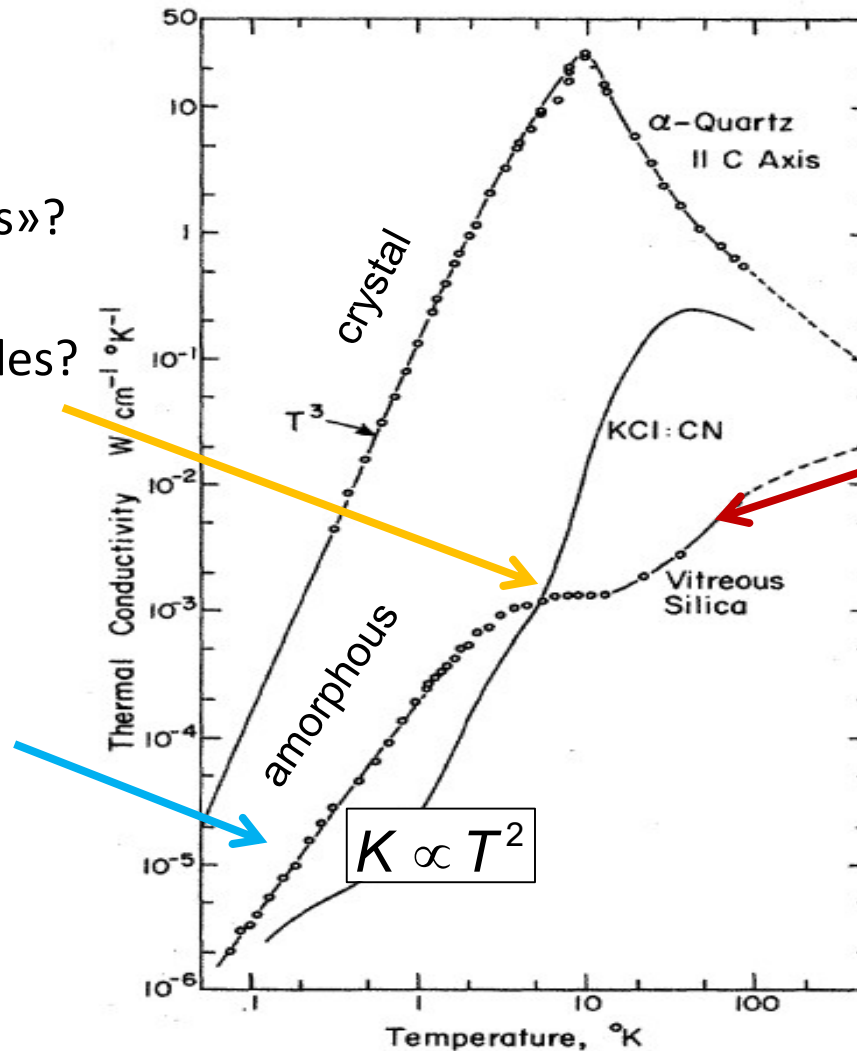
(Buchenau et al. 1992)

Strongly scattered modes?

(Wyart et al. 2010)

Resonant scattering
of phonons on
« 2-level systems »?

(Hunklinger et al 1986)



Diffusion mechanism
for heat transfer

due to
clusters vibrating
with random phases?

(Cahill et al 1988)

What is the role of **disorder** and **interfaces**
in the transportation of energy across the system?

Focus on the **vibrational** energy (phonons)

Interconnection between **Thermal** and **Mechanical** energy?

Introduction

I. Vibrations and Phonons in Amorphous Materials

II. Thermal Conductivity in Nanocomposites

Introduction

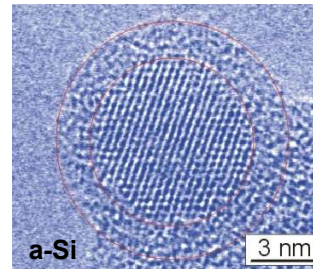
I. Vibrations and Phonons in Amorphous Materials

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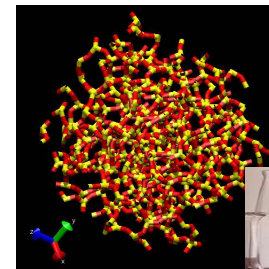
Examples of Amorphous Materials:



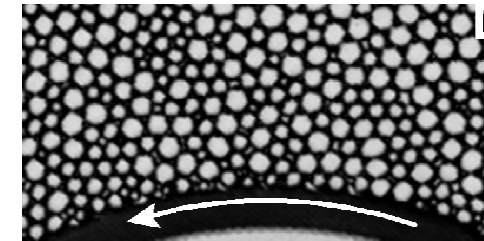
Metallic glasses, Vitreloy



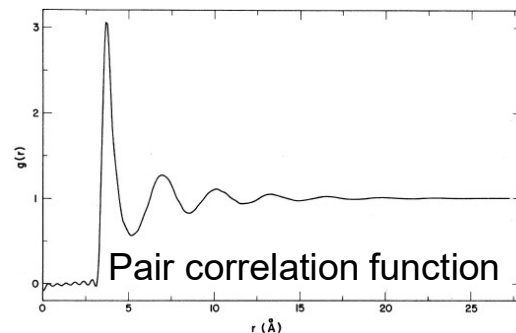
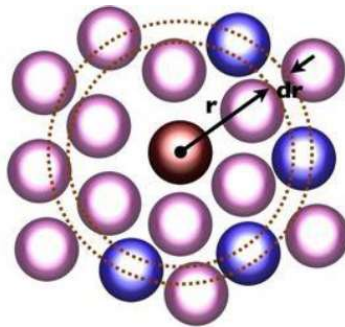
a-Si



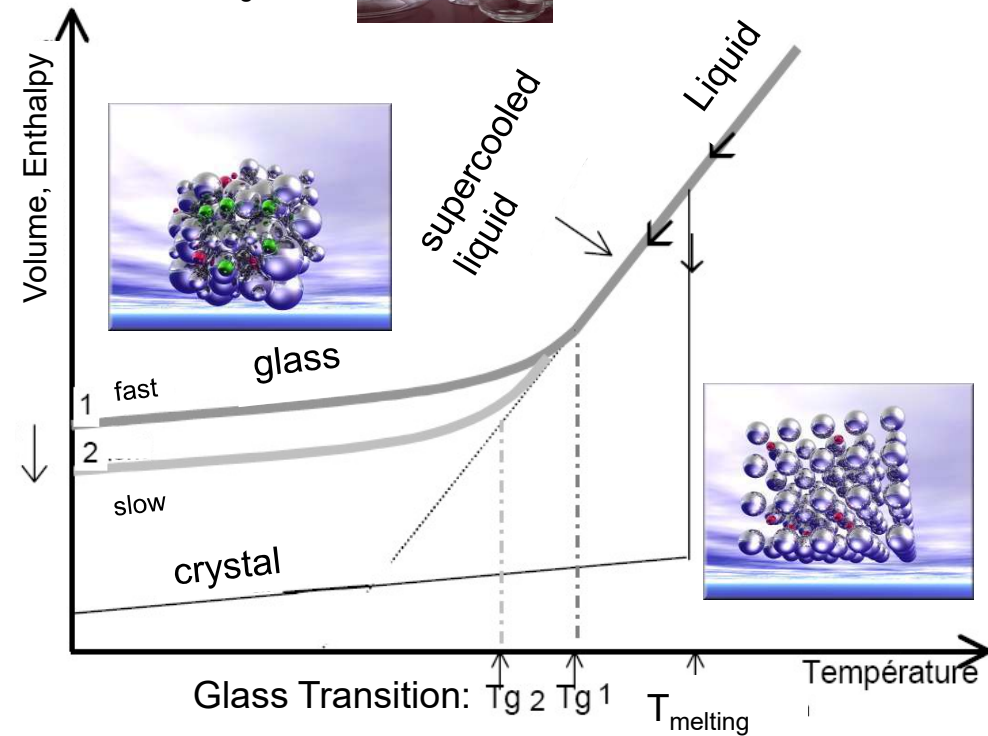
Silicate glasses



Foams



structural
relaxation ↓



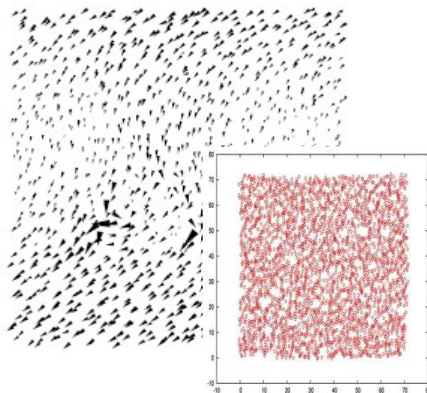
Vibration Modes in Amorphous Materials (spectral analysis)

The example of silica SiO₂ glass

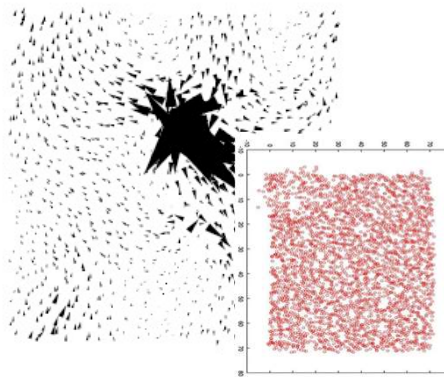
$$m_i \cdot \frac{\partial^2 u_\alpha}{\partial t^2}(\underline{r}_i, t) = - \frac{\partial E_{total}}{\partial r_{i\alpha}} \approx - \sum_j \sqrt{m_i m_j} M_{ij}^{\alpha\beta} \cdot u_\beta(\underline{r}_j, t) + f_\alpha(\underline{r}_i) \quad \text{with} \quad M_{ij}^{\alpha\beta} \equiv \frac{1}{\sqrt{m_i m_j}} \frac{\partial^2 E_{total}}{\partial r_{i\alpha} \partial r_{j\beta}}$$

$$\underline{u}(\underline{r}_i, t) \equiv \underline{r}_i(t) - \underline{r}_i^0 = \underline{U}_i / \sqrt{m_i} \cdot e^{i\omega t} \quad \text{Dynamical Matrix: Eigenvectors } \underline{U}, \text{ eigenvalues } \omega^2$$

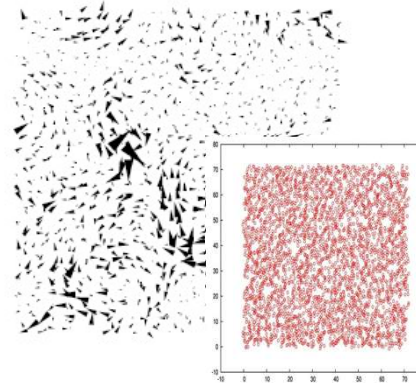
Plane Waves



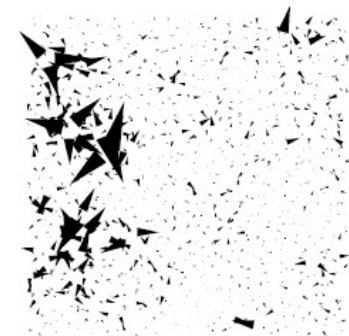
Soft Modes



Diffusons



Locons



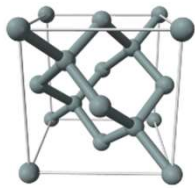
B. Mantisi et al. (2012), N. Shcheblanov et al. (2015)

P.B. Allen and
J.L. Feldman (1999)

Vibration Modes in Amorphous Materials (spectral analysis)

The example of amorphous silicon

Y. Beltukov et al. (2016)



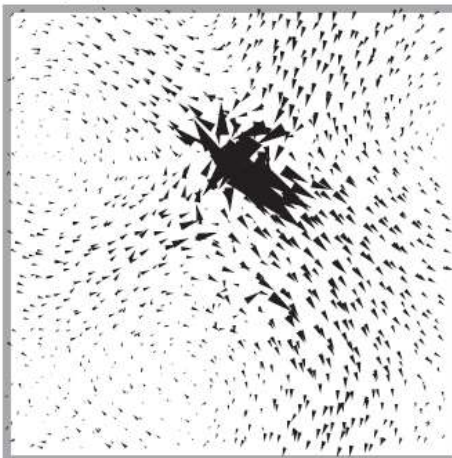
Interatomic interaction energy:

$$E_{total} = \sum_{(i,j)} f(r_{ij}) + \sum_{i,j,k} \underbrace{\Lambda \left(\cos \theta_{jik} + \frac{1}{3} \right)^2}_{\text{Bending energy}} \cdot e^{\gamma \cdot (r_{ij}-a)^{-1} + \gamma \cdot (r_{ik}-a)^{-1}}$$

Stillinger-Weber

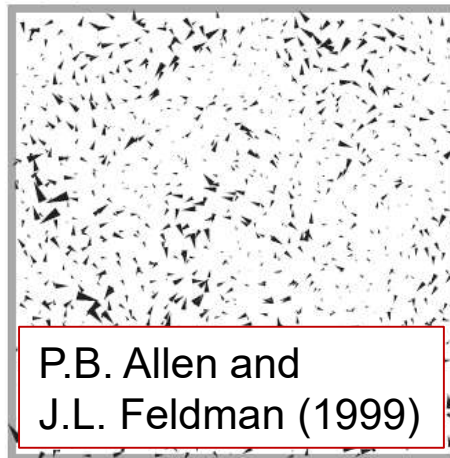
soft modes

(a) $\nu = 0.59$ THz



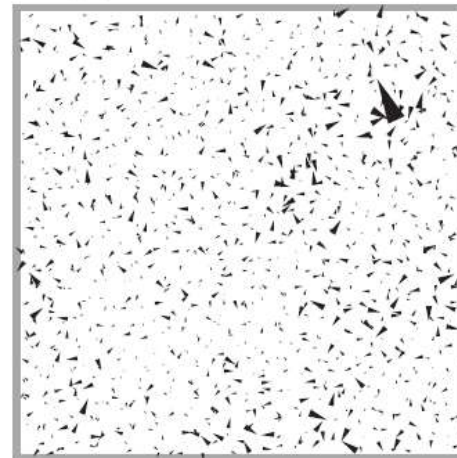
diffusons

(b) $\nu = 2.67$ THz



P.B. Allen and
J.L. Feldman (1999)

(c) $\nu = 9.95$ THz



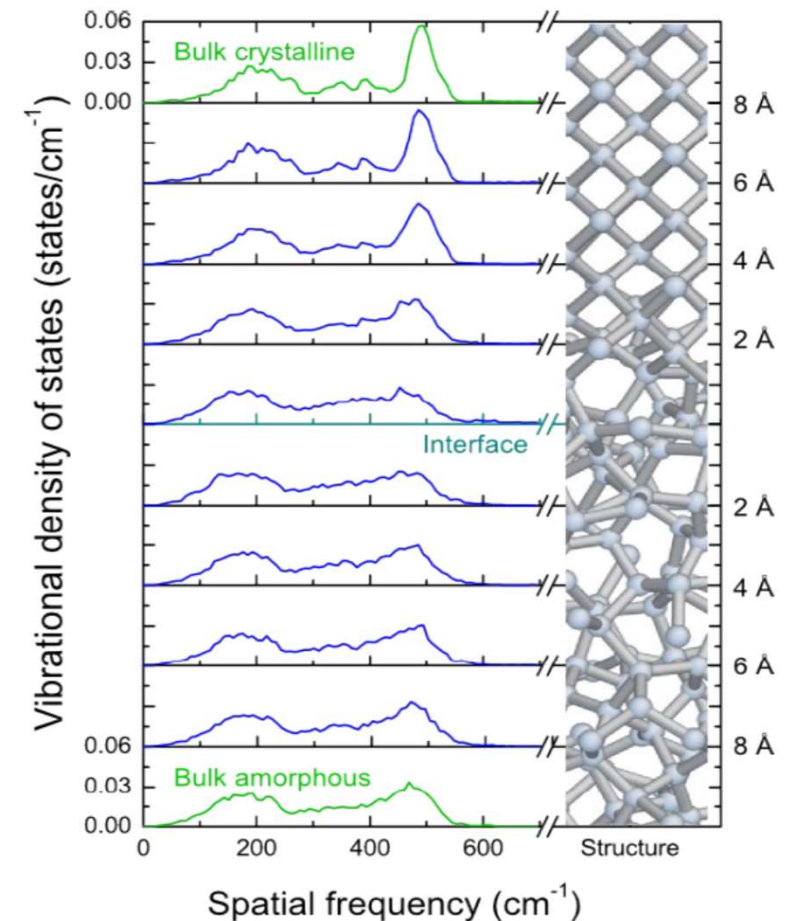
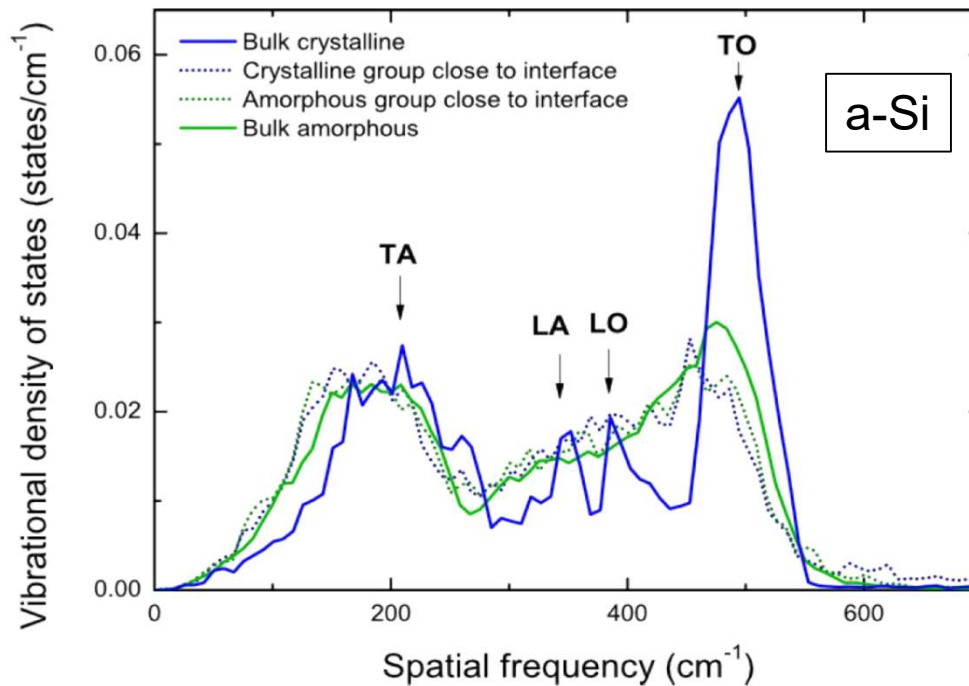
locons

(d) $\nu = 17.39$ THz



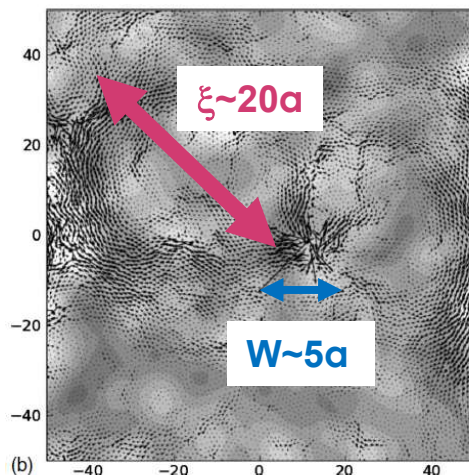
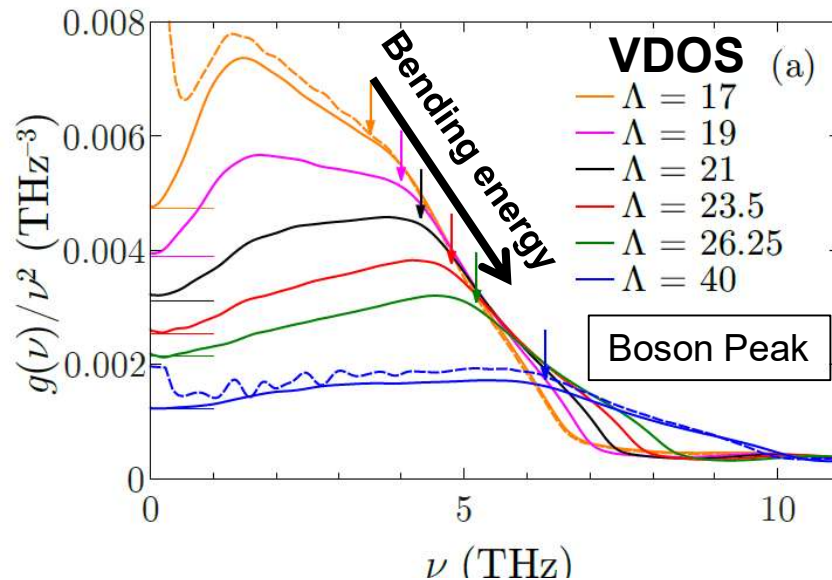
Vibrational Density of States

Structural sensitivity, from crystal to amorphous solid



A. France-Lanord et al. (2014)

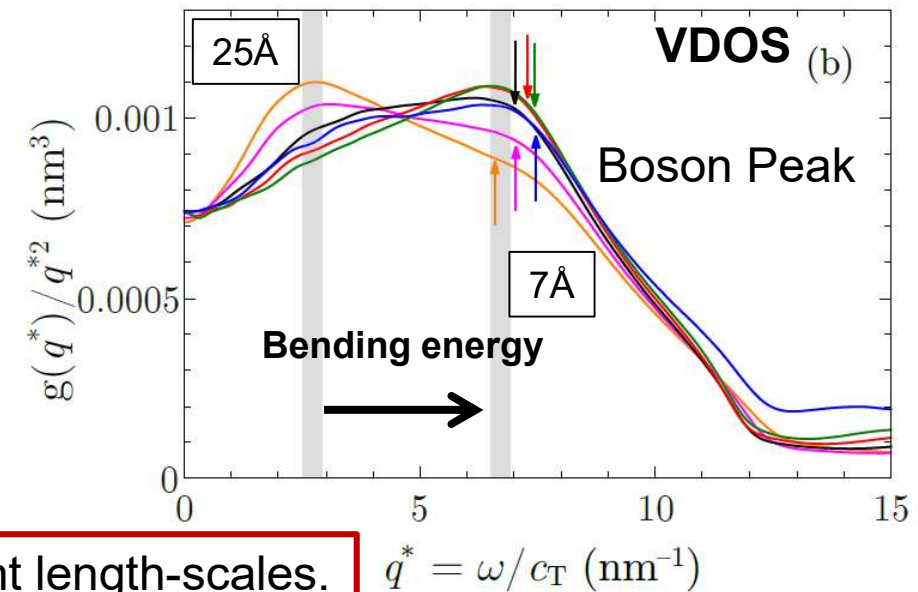
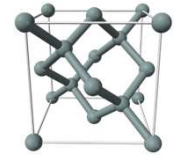
Normalized Vibrational Density of States in a-Si



Interatomic interactions:

$$E_{SW} = \sum_{(i,j)} f(r_{ij}) + \sum_{i,j,k} \Lambda \left(\cos \theta_{jik} + \frac{1}{3} \right)^2 e^{\gamma \cdot (r_{ij}-a)^{-1} + \gamma \cdot (r_{ik}-a)^{-1}}$$

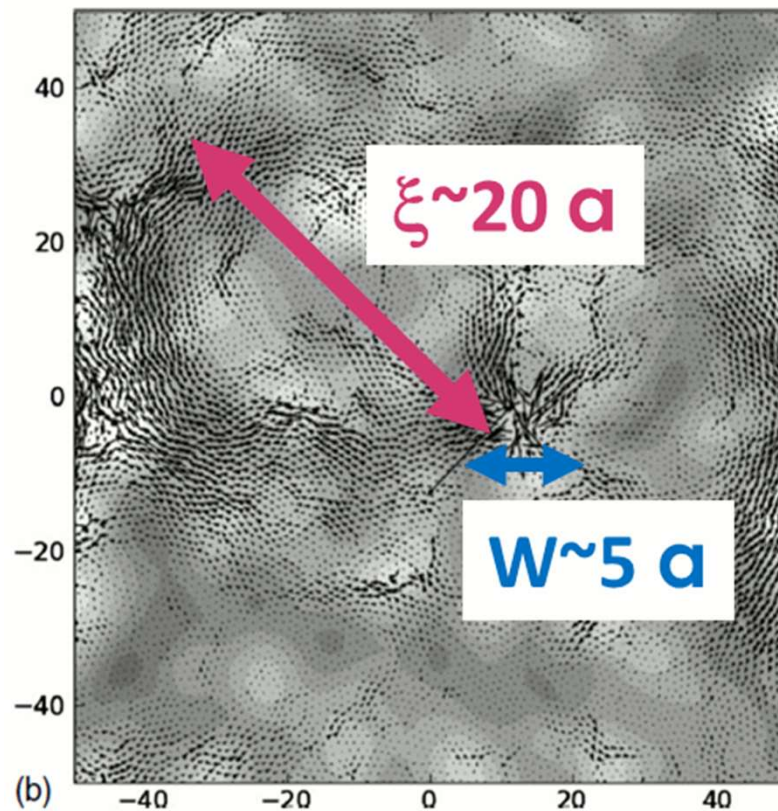
Λ , bending rigidity



Two different length-scales.
Waves scattering over
correlated structures.

Y. Beltukov et al. (2016)
A. Tanguy (2015)

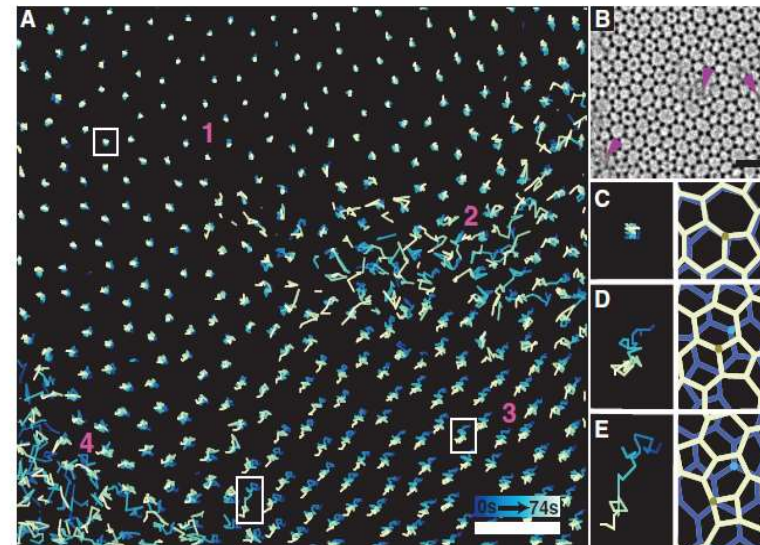
Characteristic Length-scales in Amorphous Solids:



M. Tsamados et al. (2009)

Imaging Atomic Rearrangements in Two-Dimensional Silica Glass: Watching Silica's Dance

Pinshane Y. Huang,¹ Simon Kurasch,^{2*} Jonathan S. Alden,^{1*} Ashvini Shekhawat,³
Alexander A. Alemi,³ Paul L. McEuen,^{3,4} James P. Sethna,³ Ute Kaiser,² David A. Muller^{1,4†}

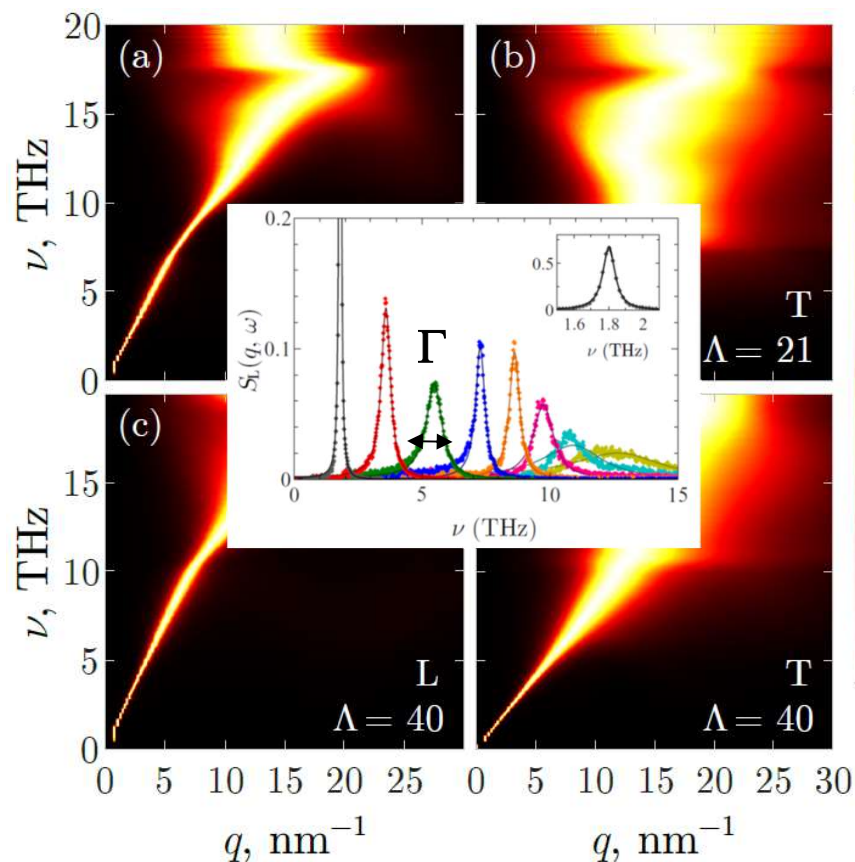


P.Y. Huang et al. Science (October 2013)

Dynamical Structure Factor in a-Si:

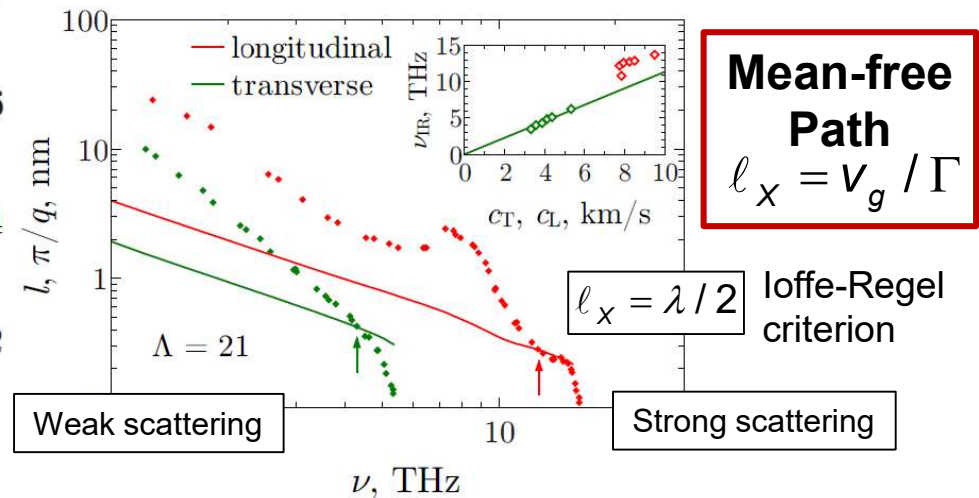
$$S_L(q, \omega) = \int e^{-i\omega t} \langle \rho(\bar{q}, t) \cdot \rho(-\bar{q}, 0) \rangle dt$$

Y. Beltukov et al. (2016)



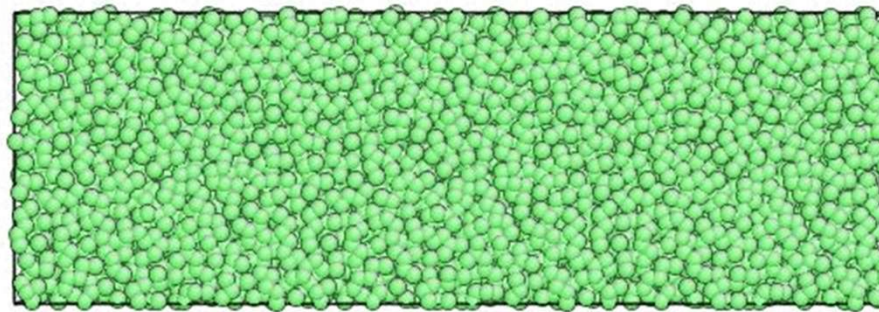
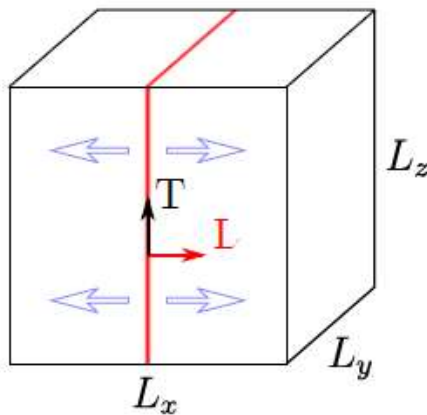
Damped Harmonic Osc fit:

$$S_{L,T}(q, \omega) = \frac{A}{(\omega^2 - \omega_{L,T}^2(q))^2 + \omega^2 \Gamma^2}$$

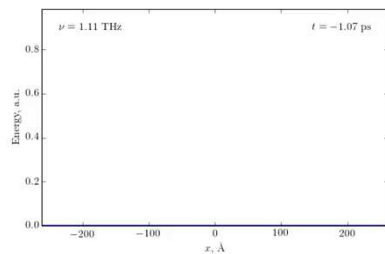


Scattering of Transverse
and then Longitudinal waves

Wave-Packets Dynamics in a-Si:

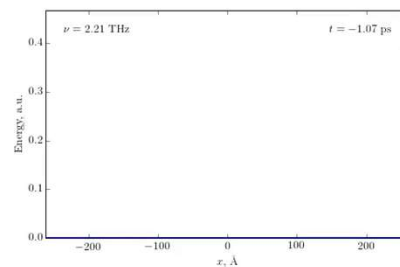


$$f_i^{\text{exc}}(\omega, t) = f_\eta \exp \left(i\omega t - \frac{t^2}{2\tau_{\text{exc}}^2} - \frac{x_i^2}{2w^2} \right)$$

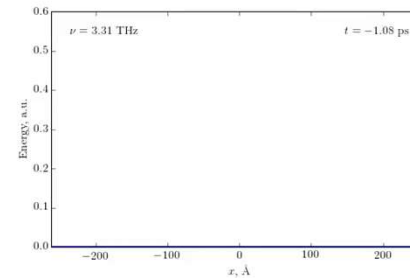


Ballistic Regime
Propagons

$$\omega \ll \omega_{\text{IR}}$$

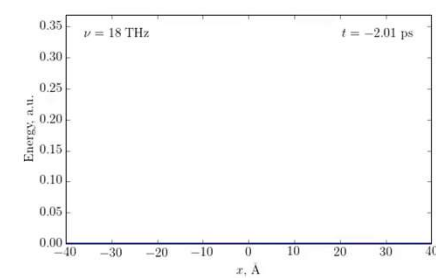


Mixed Regime
(Boson Peak)



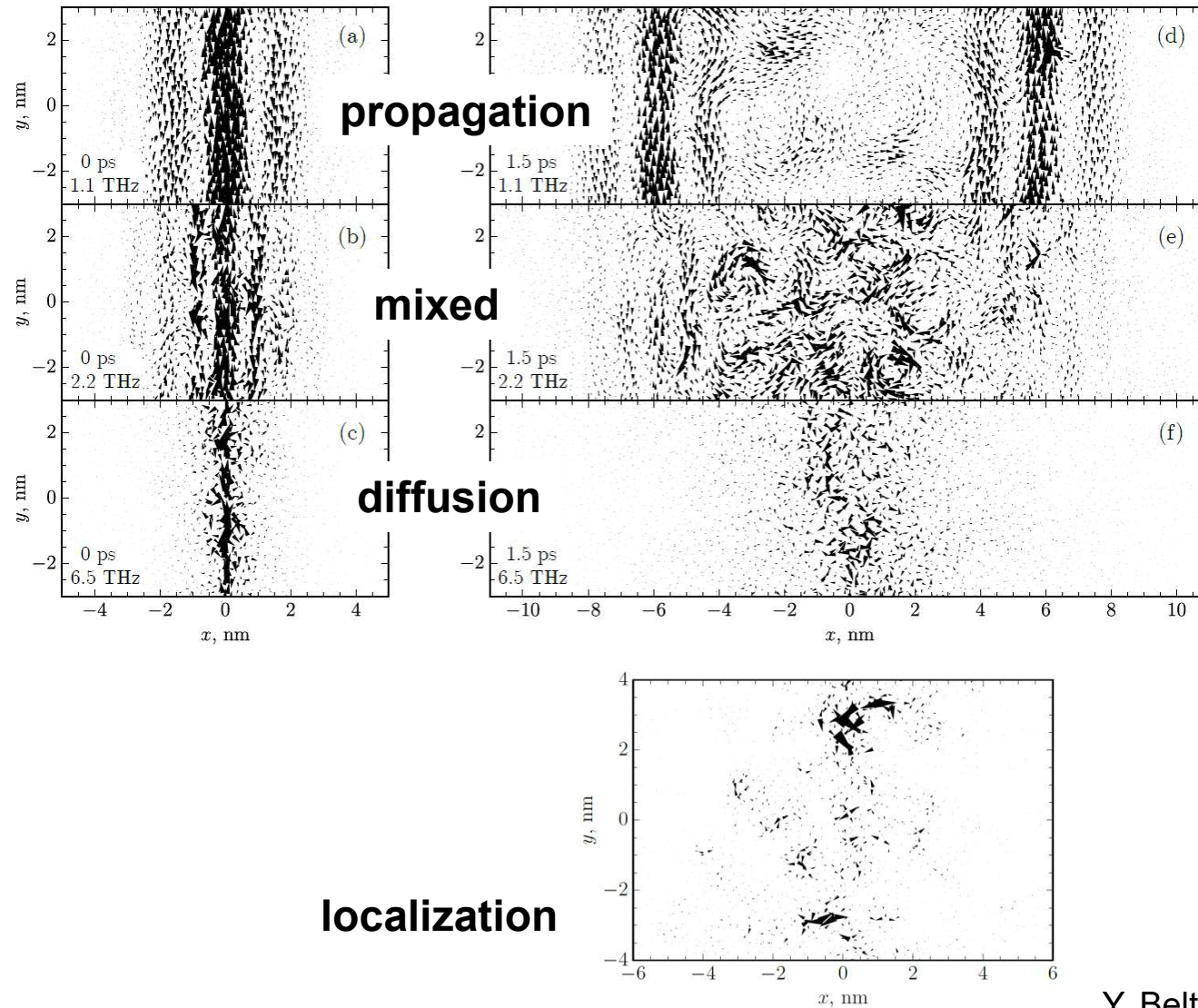
Diffusive Regime
Diffusons

$$\omega > \omega_{\text{IR}}$$

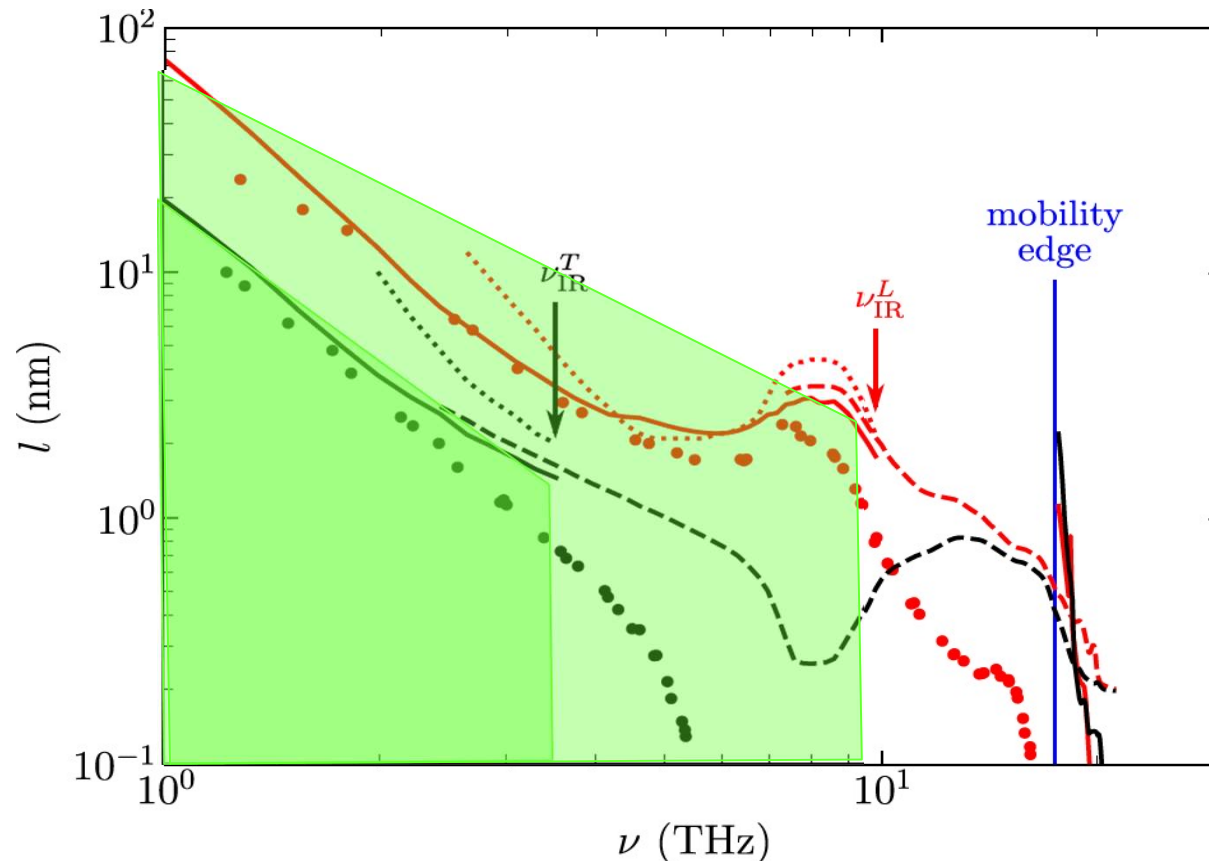


Localisation
Locons

Y. Beltukov et al. (2018)



Attenuation Length Comparison to the Dynamical Structure Factor



--- Attenuation length
 $P_{\text{diff-prop}}(l) = P_{\text{diff-prop}}(0) / e$

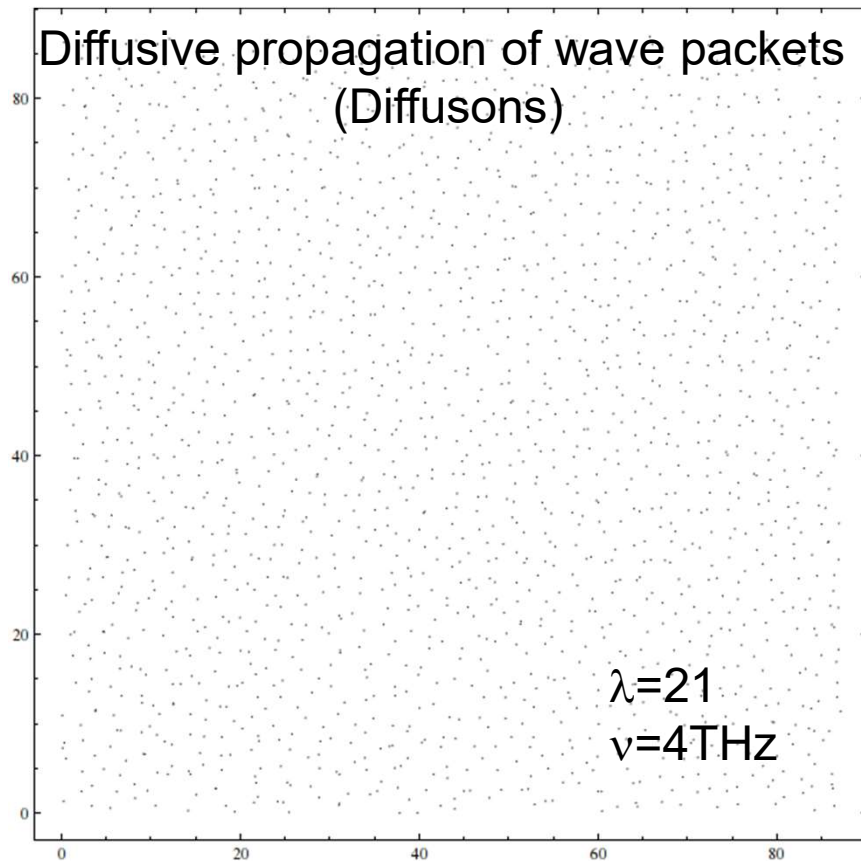
— Mean-free path
(Beer-Lambert)

● Dyn. Struct. Factor
 $l = c_{L,T} / \Gamma$

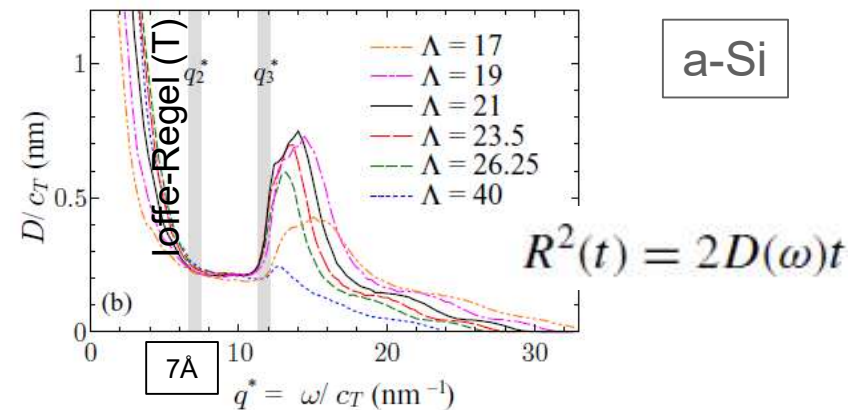
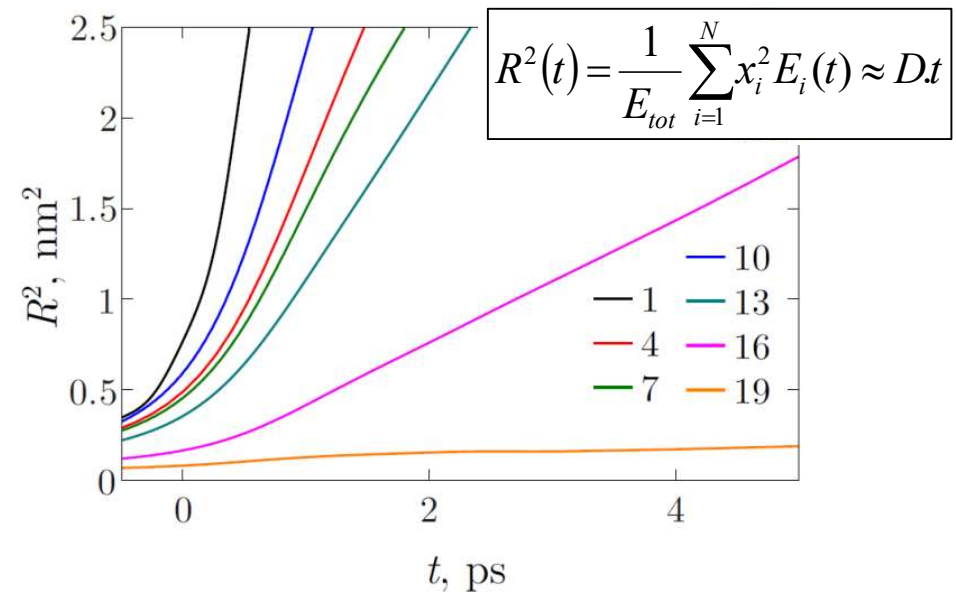
Transportation of
Energy even **beyond**
the Ioffe-Regel criterion

The Mean-free path (and
from the DSF) is defined
only in the Ballistic regime

Determination of **Diffusivity** from **Wave-Packets** motion



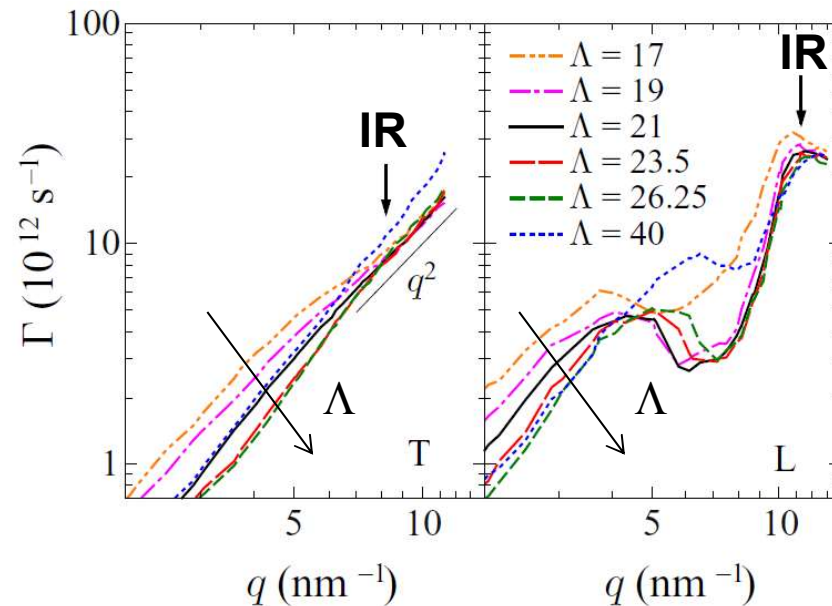
Minimum **diffusivity**
at the Ioffe-Regel criterion
of Transverse waves.



Allen and Feldman (1999)

Ballistic Regime

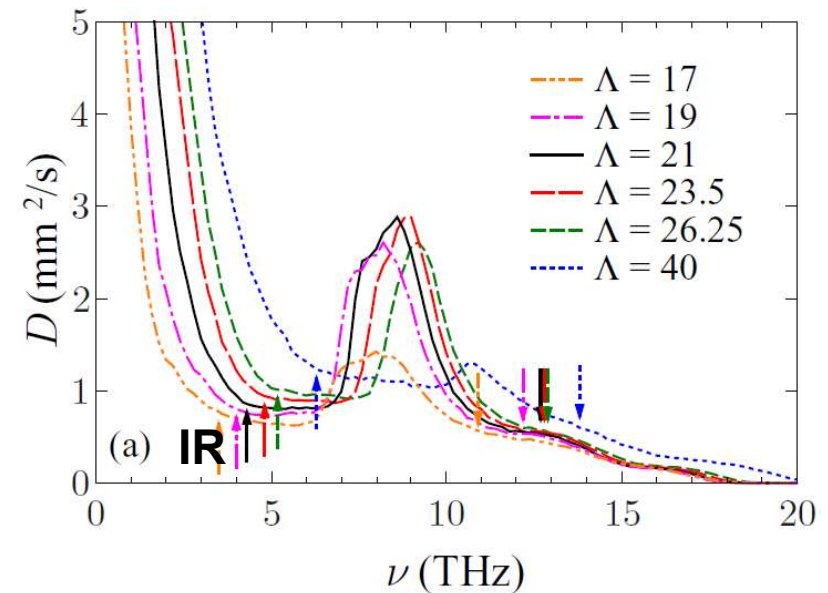
Attenuation for $\omega < \omega_{IR}$



$$\kappa^{prop}(T) = \frac{1}{3} \int_0^{\omega_{IR}} C(\omega, T) v_g^2(\omega) \Gamma^{-1}(\omega) g(\omega) d\omega$$

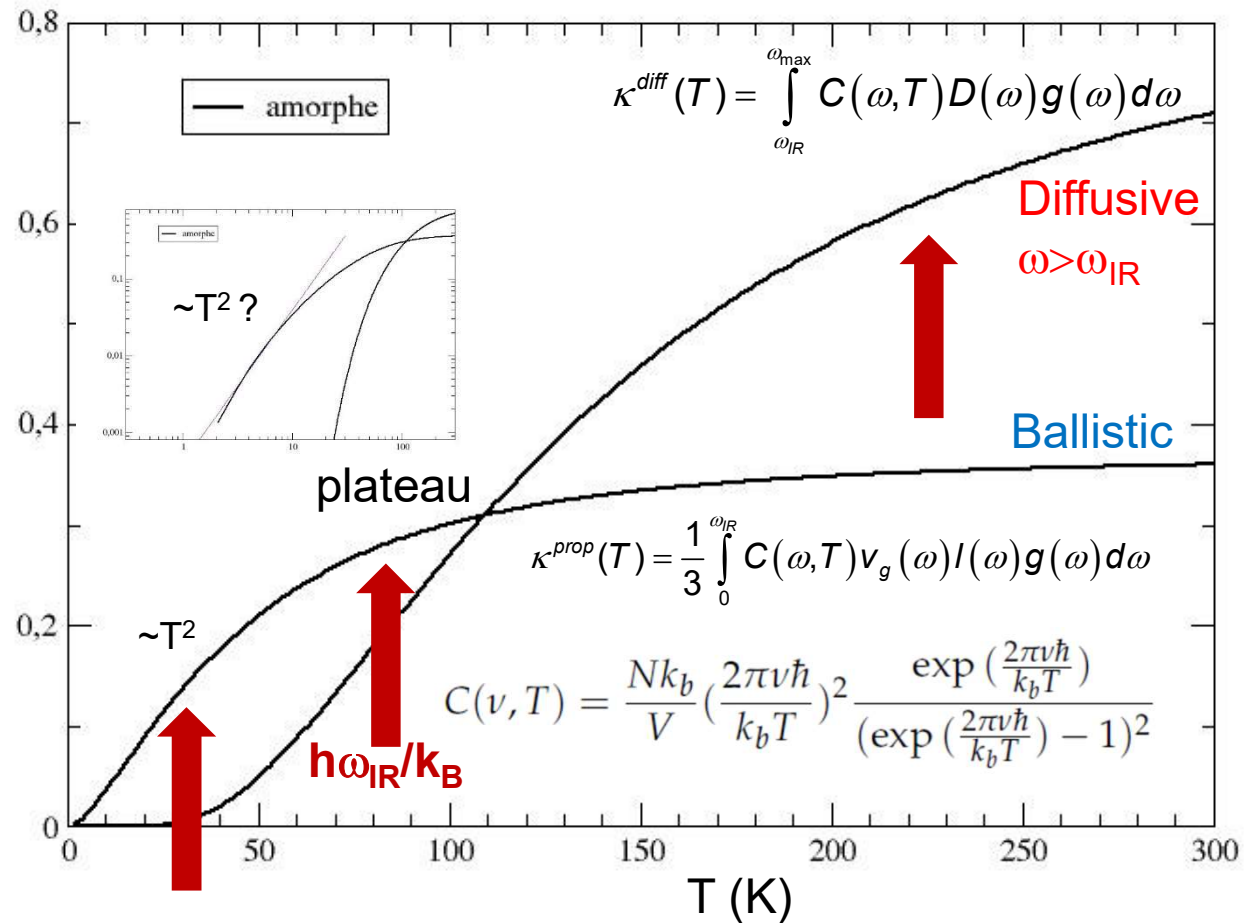
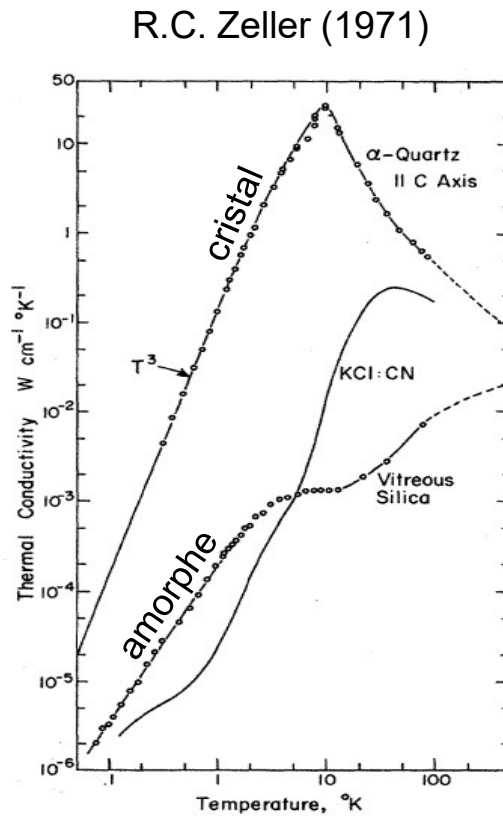
Diffusive Regime

Diffusivity for $\omega \geq \omega_{IR}$



$$\kappa^{diff}(T) = \int_{\omega_{IR}}^{\omega_{max}} C(\omega, T) D(\omega) g(\omega) d\omega$$

Thermal Conductivity

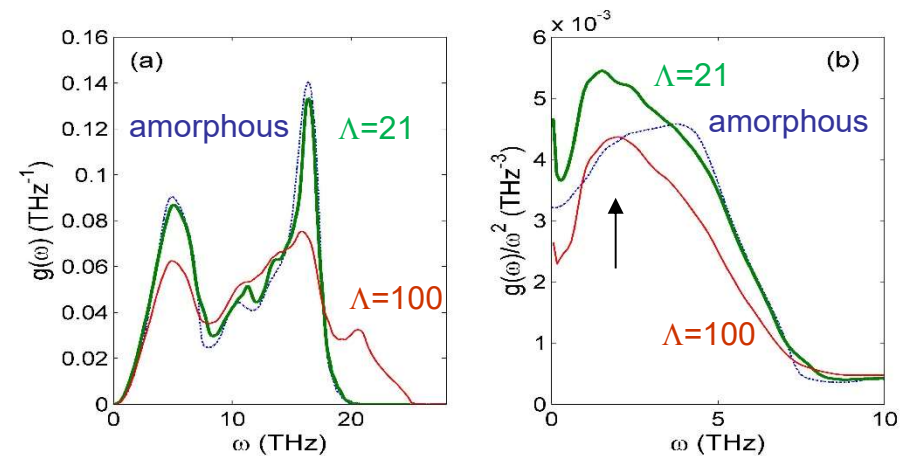
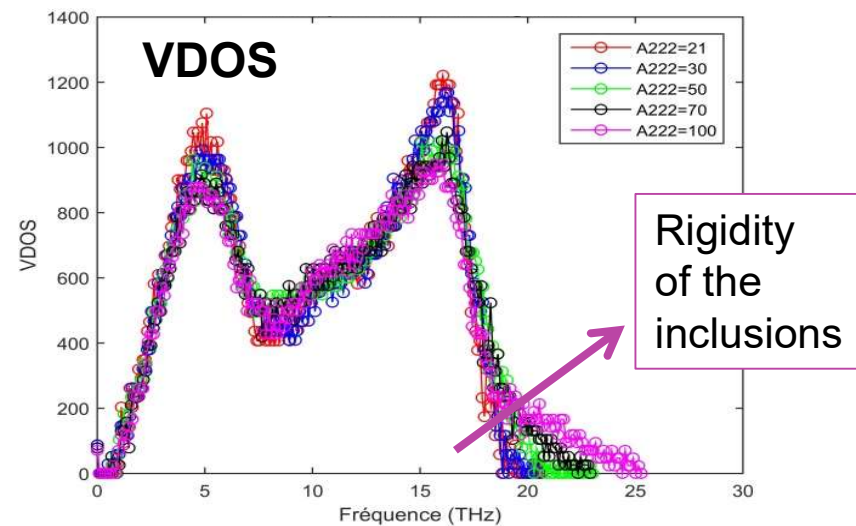
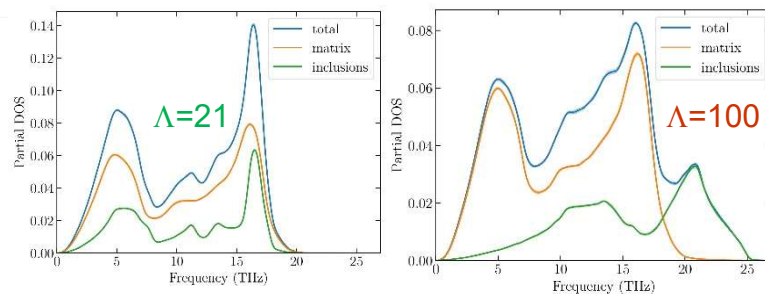
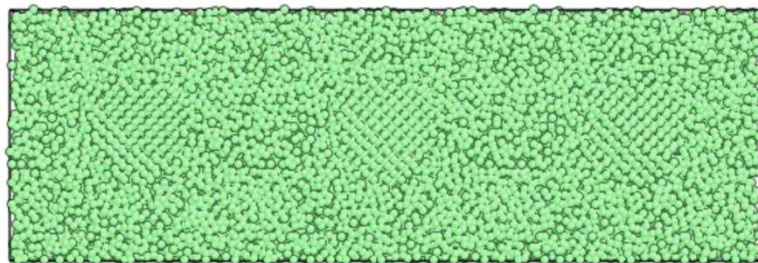
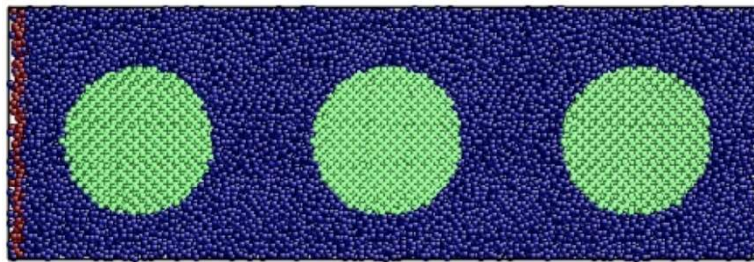


Introduction

I. Vibrations and Phonons in Amorphous Materials

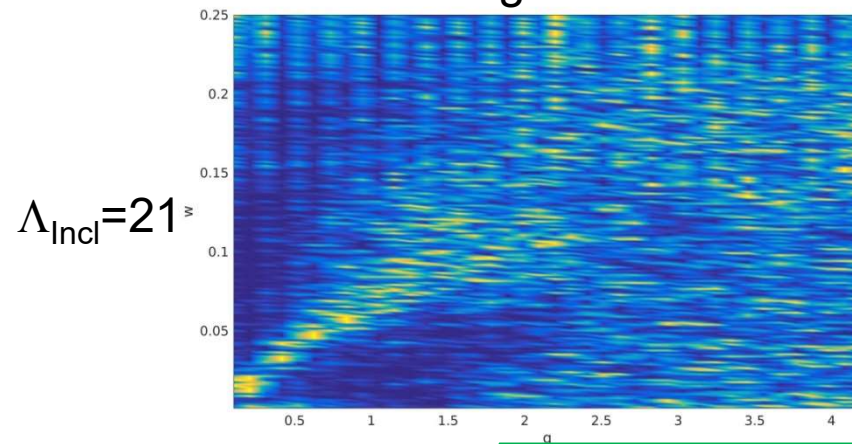
II. Thermal Conductivity in Nanocomposites

Molecular Dynamics Simulations of the nanocomposite a-Si/c-Si

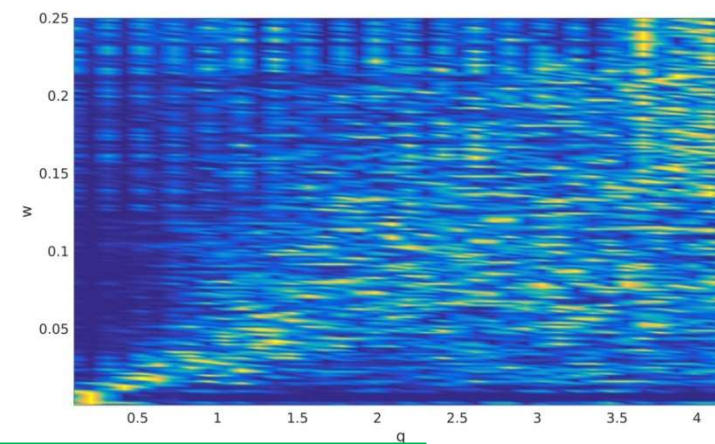


Dynamical Structure Factor in a-Si/c-Si

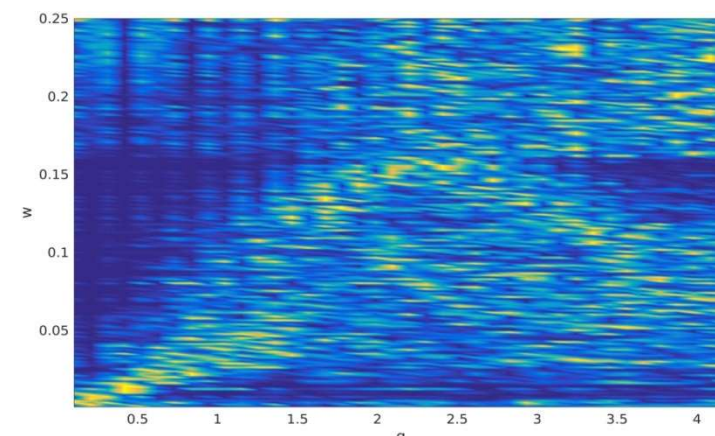
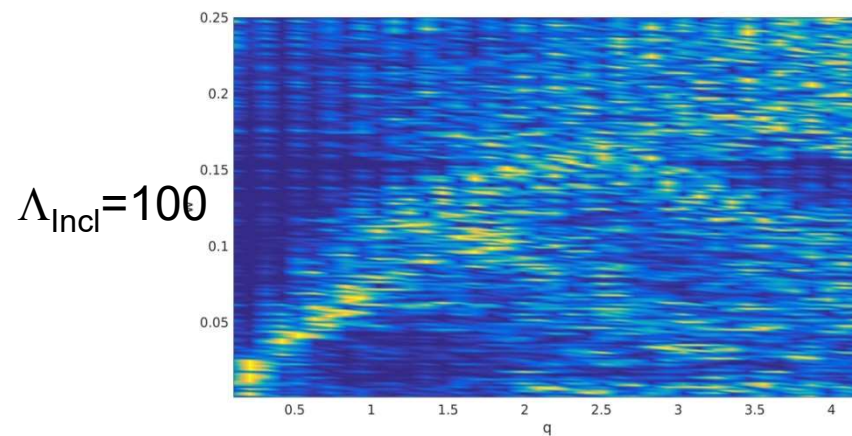
Longitudinal



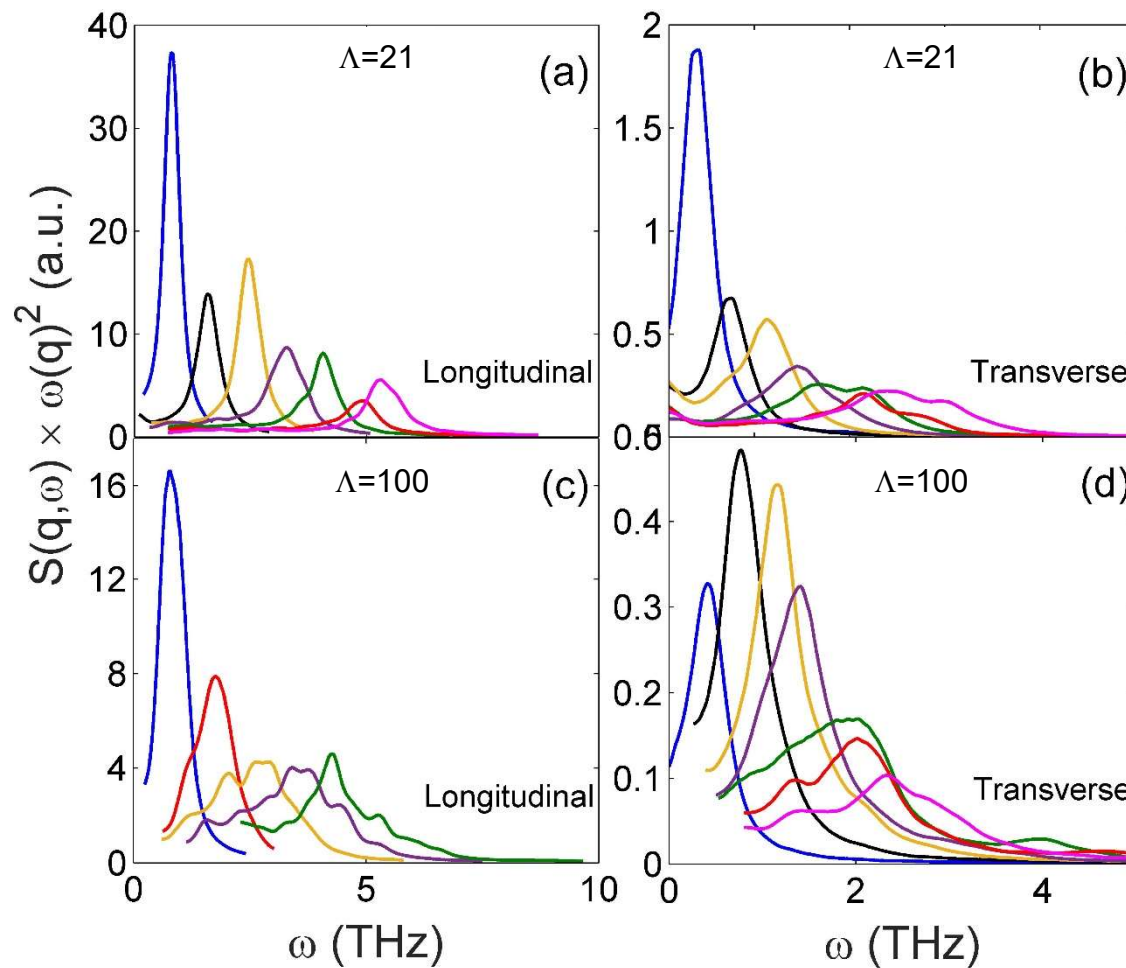
Transverse



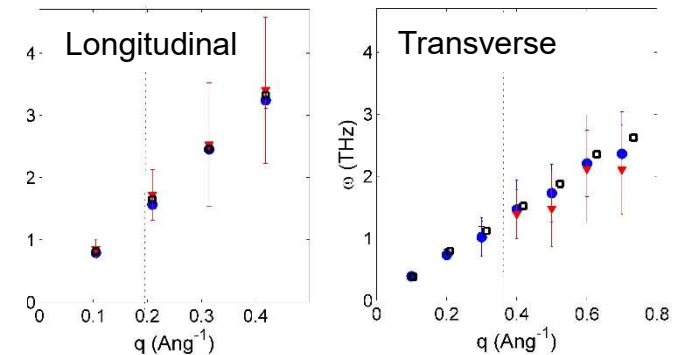
Splitting in the Dynamical Structure Factor



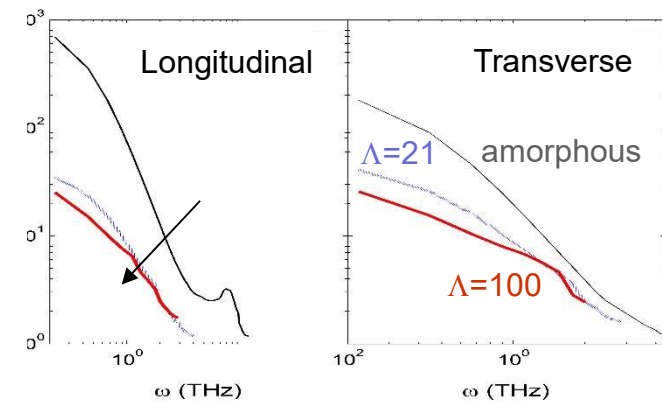
Dynamical Structure Factor in a-Si/c-Si



Dispersion Relation:

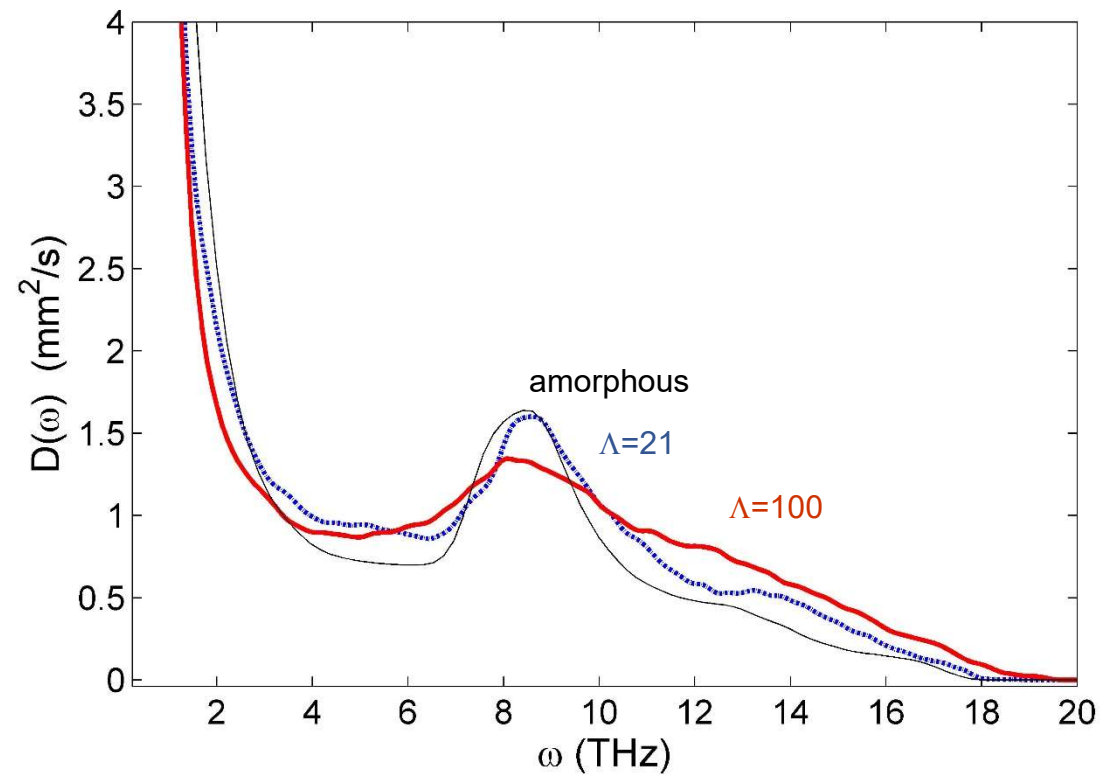


Mean-Free Path:



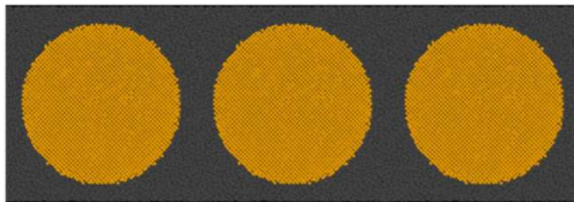
Anticipation of the
Ioffe-Regel cross-over

Comparable Diffusivity of Phonons in a-Si/c-Si



Diffusivity dominated by the amorphous matrix

Thermal Conductivity in a-Si / c-Si nanocomposite

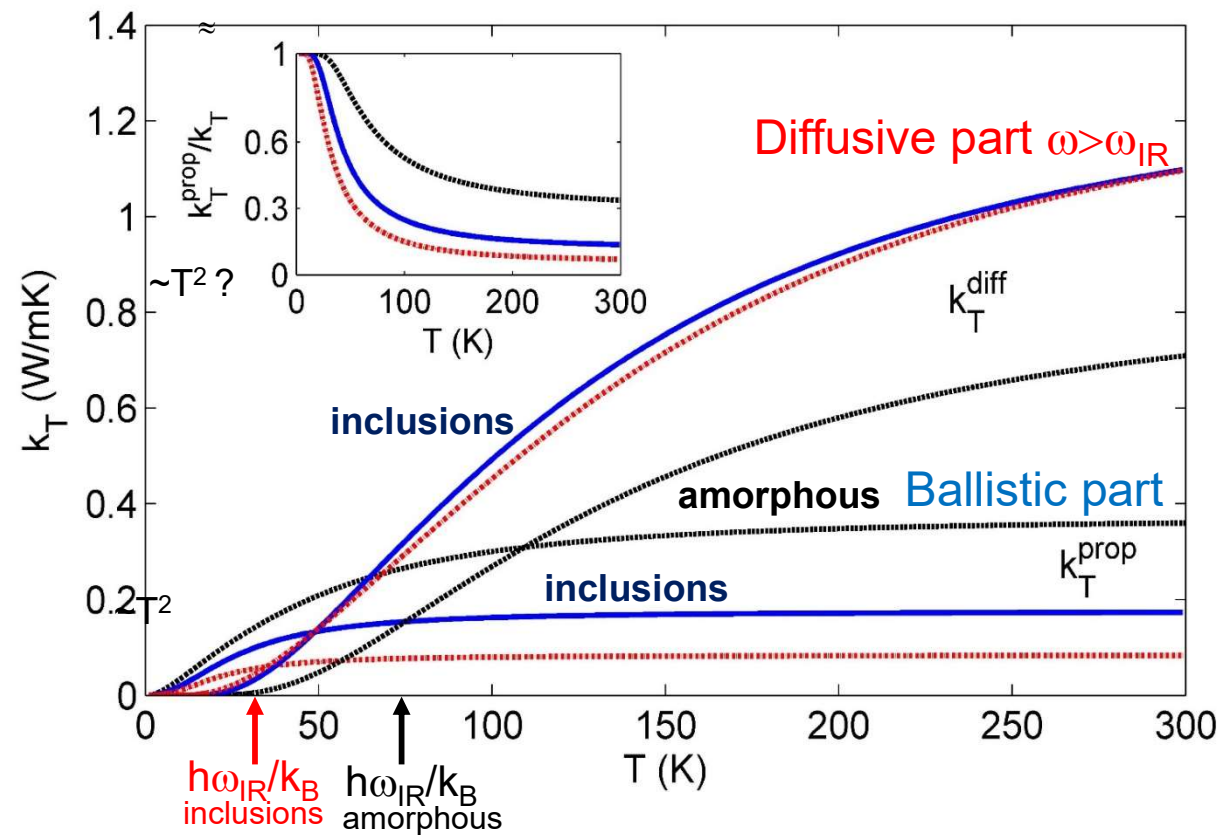


Rigidity contrast:

— $E_i / E_m \approx 1$

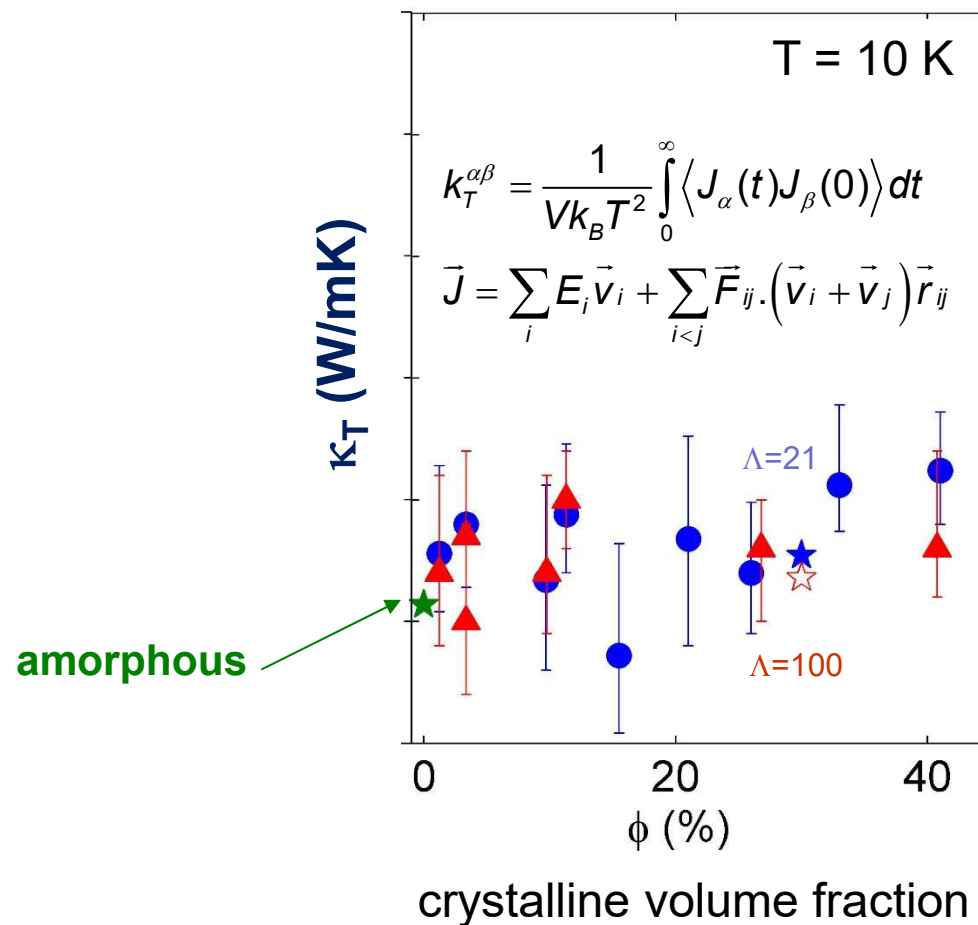
— $E_i / E_m \approx 4.6$

The transition frequency
to multiple scattering ω_{IR}
is smaller ↘



Enhancement of the diffusive contribution to Thermal conductivity in the nanocomposite

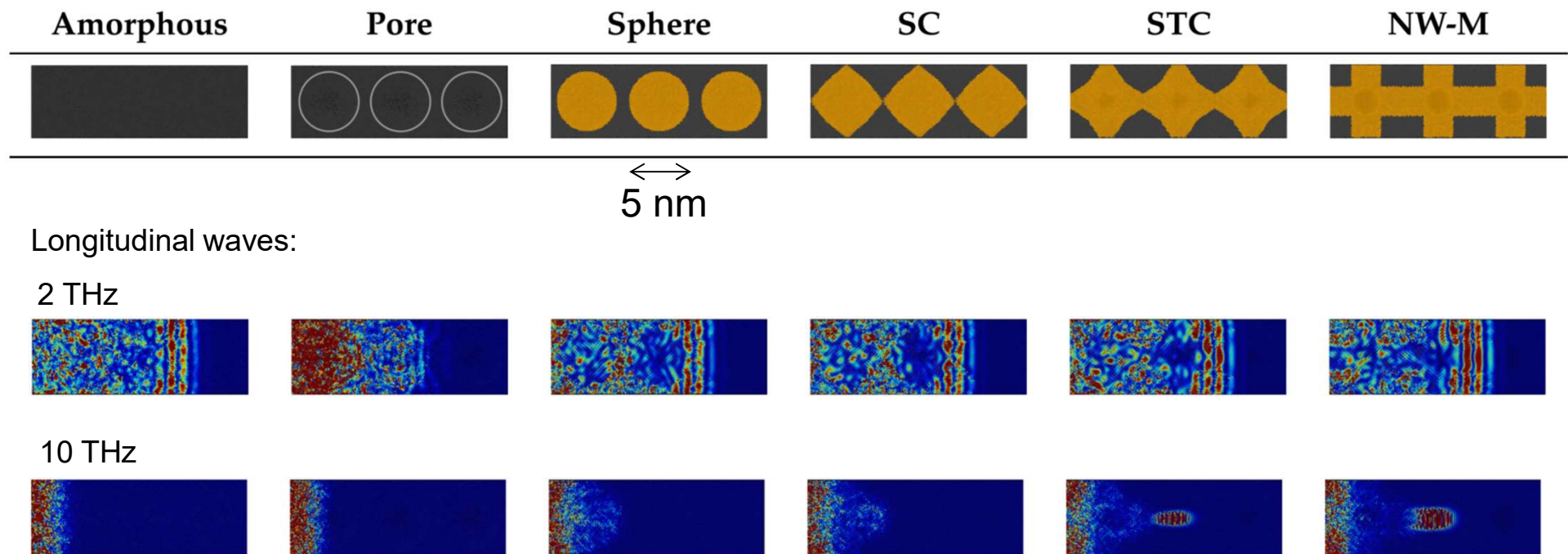
Green-Kubo calculation of Thermal Conductivity



No clear effect
of the Nano-inclusions
on the Thermal Conductivity
in the low T regime

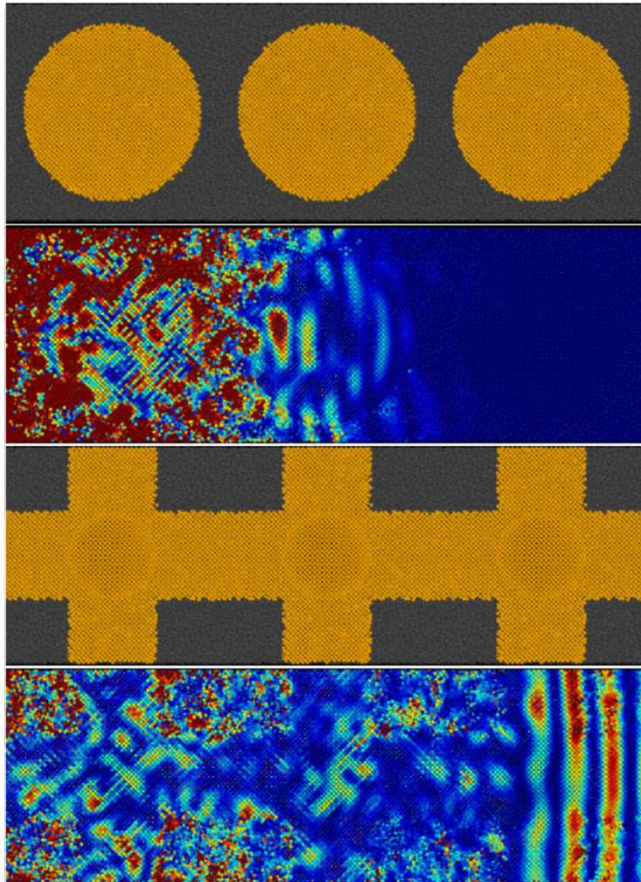
Thermal Conductivity in a-Si / c-Si nanocomposite

Role of the interconnections between the inclusions

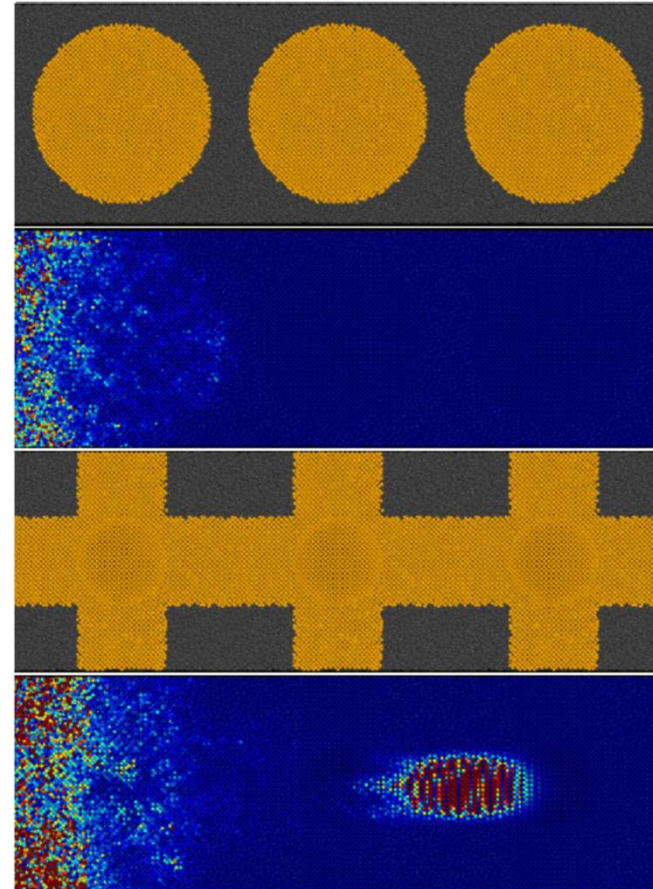


Wave Packets propagation in case of percolating vs. Non-percolating nanostructure

2 THz

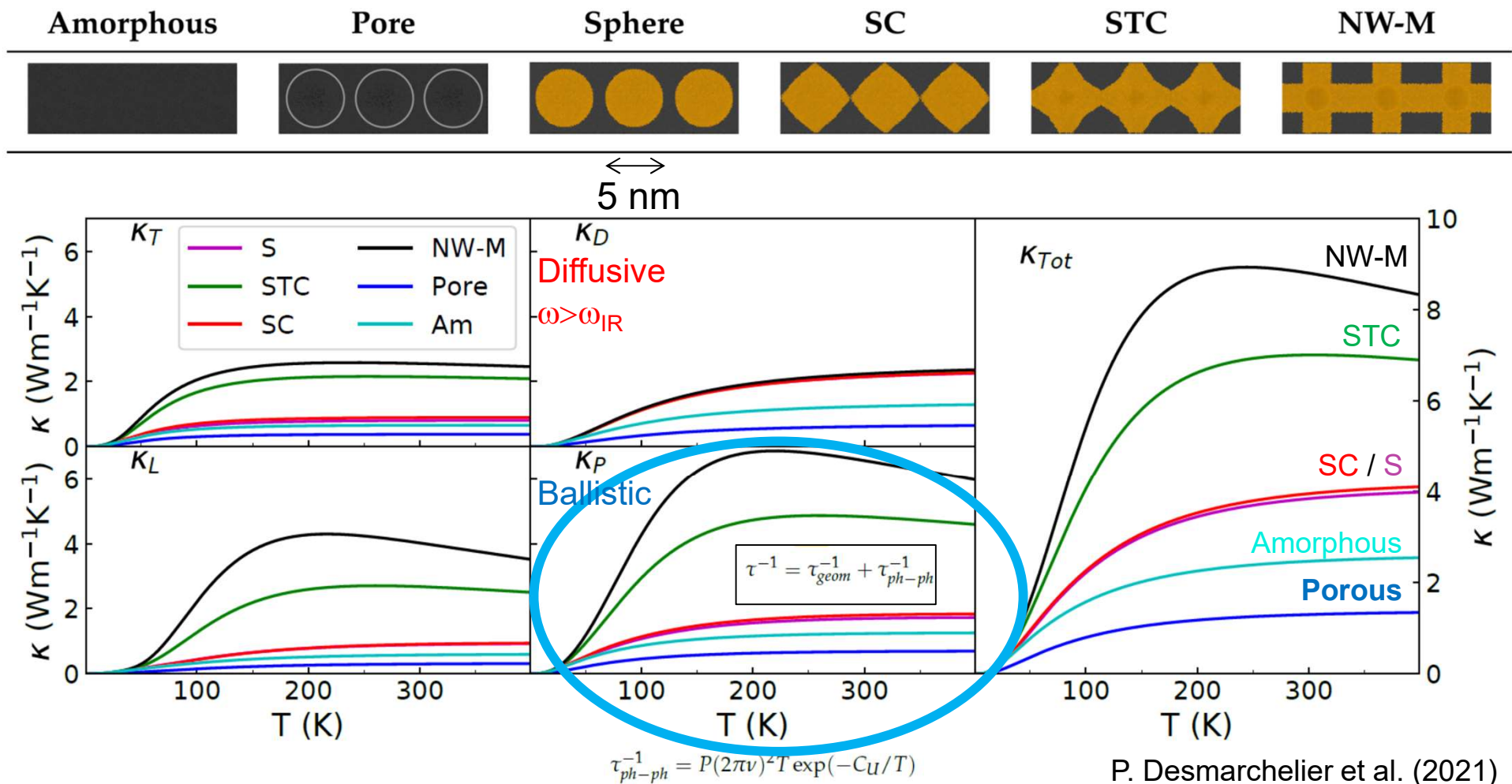


10 THz



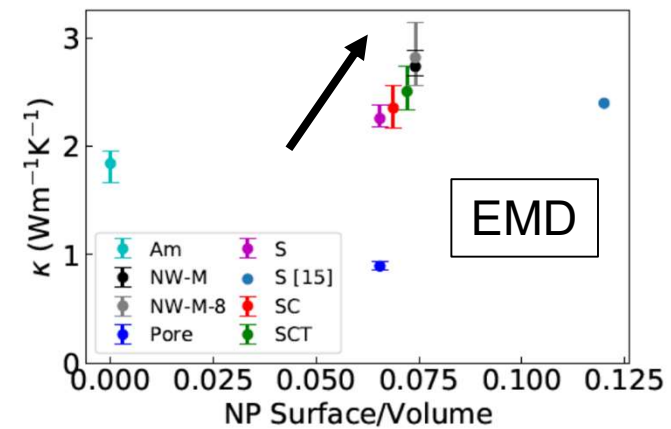
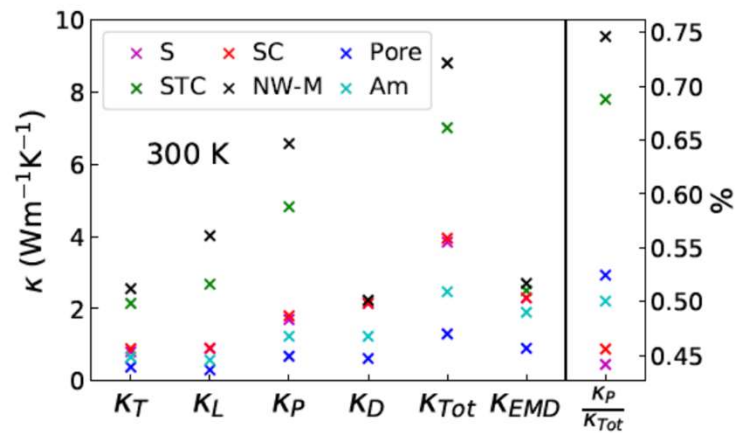
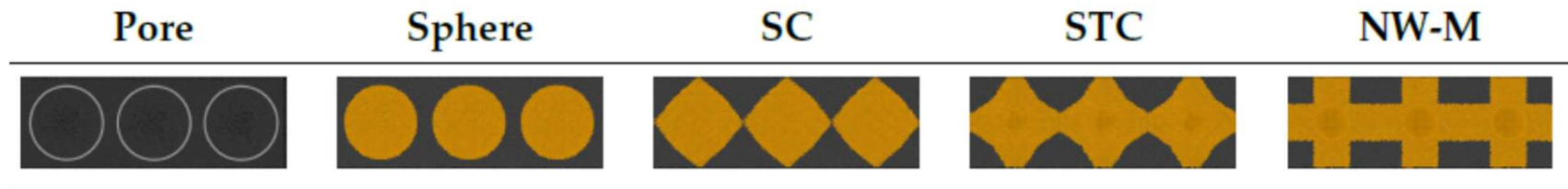
P. Demarchelier et al. (2020)

Thermal Conductivity in a-Si / c-Si nanocomposite



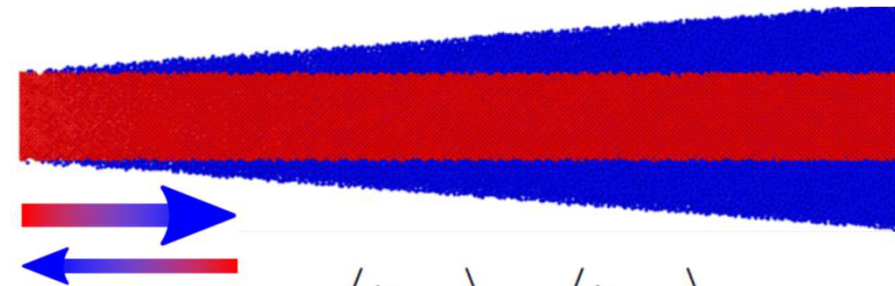
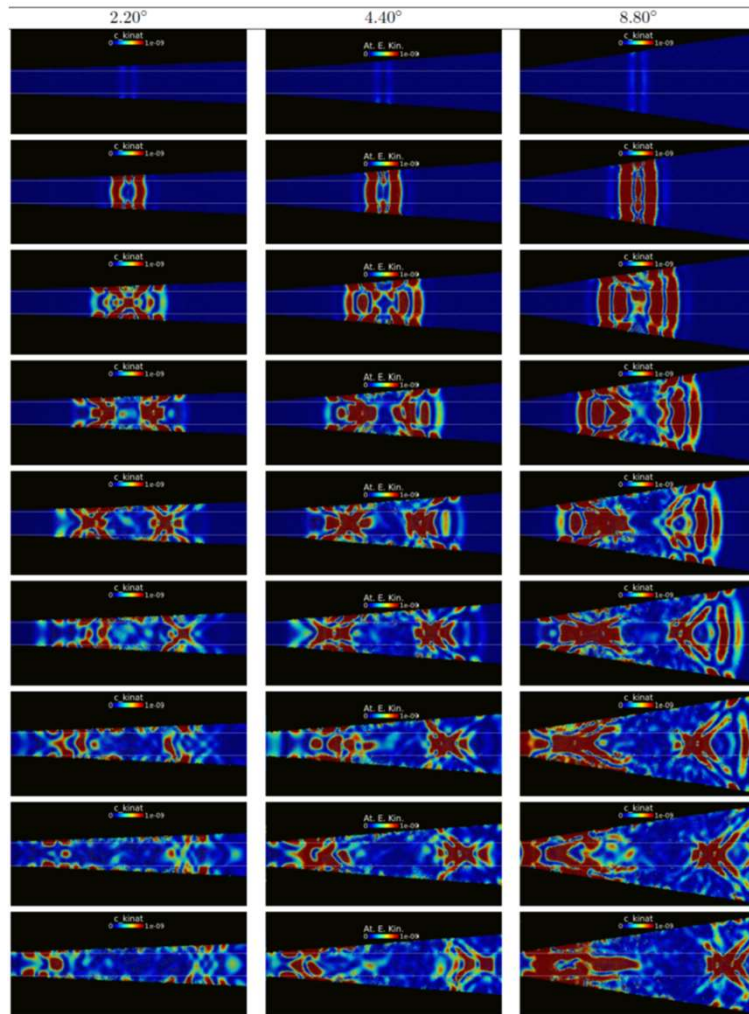
Thermal Conductivity in a-Si / c-Si nanocomposite

Role of the interconnections between the inclusions

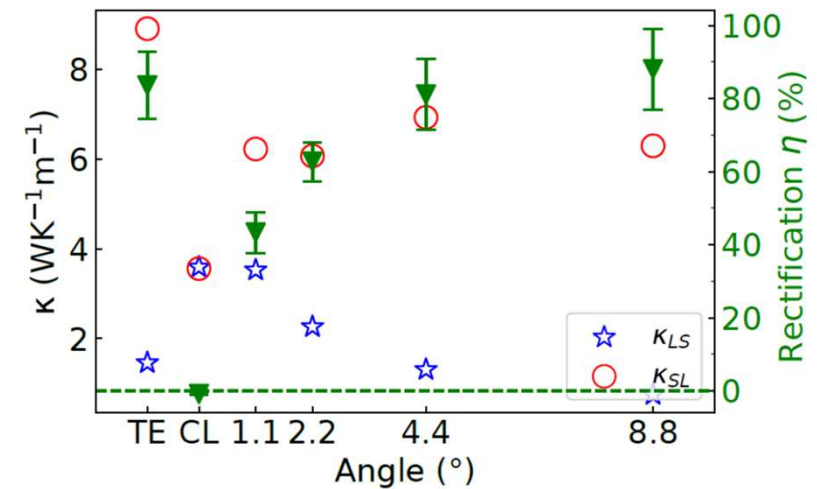


Thermal Rectification in c-Si nanowire with amorphous shell

2 THz



$$\eta = \frac{\langle \kappa_{SL} \rangle - \langle \kappa_{LS} \rangle}{\langle \kappa_{SL} \rangle}$$



Conclusion

Phononic contribution to Thermal Conductivity in Amorphous Solids:

Low scattering, strong scattering and **localization** of wave-packets in amorphous materials:

- the **Ioffe-Regel** cross-over is the transition between low and strong scattering
- the **Mobility Edge** is the transition to localization.

Propagative contribution to Thermal Conductivity dominates at low frequencies. Above the Boson Peak frequency, the **diffusive contribution** dominates.

In Nanocomposites:

The presence of **nanoinclusions** changes the **ratio** between the propagative and the diffusive contribution to thermal conductivity.

The **interconnection** (percolation) between the inclusions has a crucial effect.

Playing with **interfaces** paves the way for **thermal management** (rectification)

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Thank You